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Pelka et al.

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(54) **ILLUMINATION LENS FOR SHORT-THROW LIGHTING**

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F21W 131/305 (2006.01)
F21W 131/405 (2006.01)
F21Y 101/02 (2006.01)

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F25D 27/00 (2013.01); **G02B 19/0028**
(2013.01); **G02B 19/0061** (2013.01); **F21W**
2131/305 (2013.01); **F21W 2131/405**
(2013.01); **F21Y 2101/02** (2013.01)

USPC **362/311.02**; 362/311.05; 362/311.1

(58) **Field of Classification Search**

USPC 362/235, 311.02, 311.05, 311.1
See application file for complete search history.

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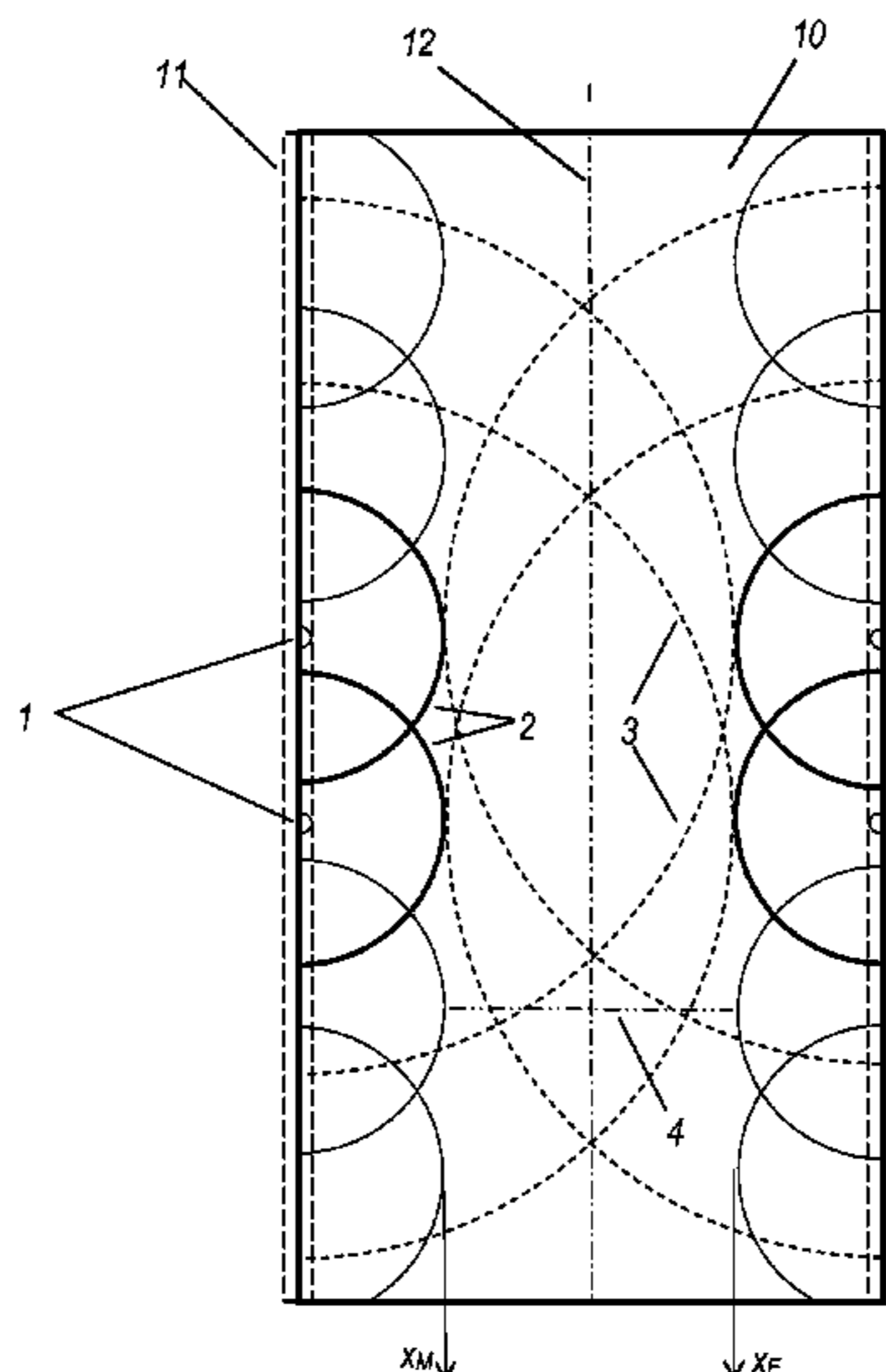
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PLC

(57) **ABSTRACT**

Exemplary embodiments of the present invention relate to
wide-angle illumination patterns for short-throw lighting. An
illumination lens according to an exemplary embodiment
includes a light incident surface that defines a cavity and has
a flat top and lateral flanks, the light incident surface being
configured to receive light from an underlying light emitting
element. The illumination lens also includes a light exiting
surface having a central indentation and a surrounding toroid,
wherein the flat top faces the central indentation.

10 Claims, 11 Drawing Sheets



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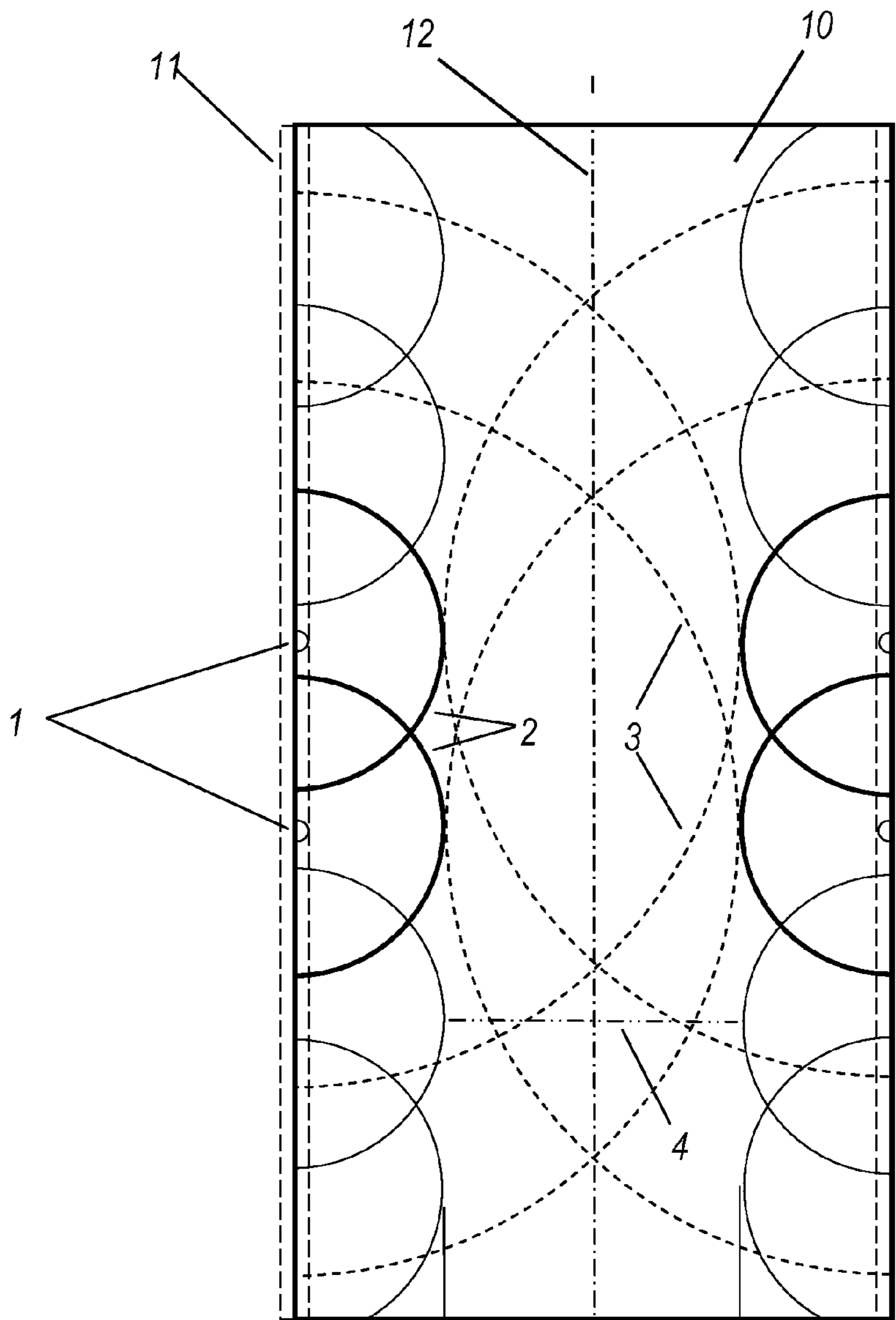


FIG 1

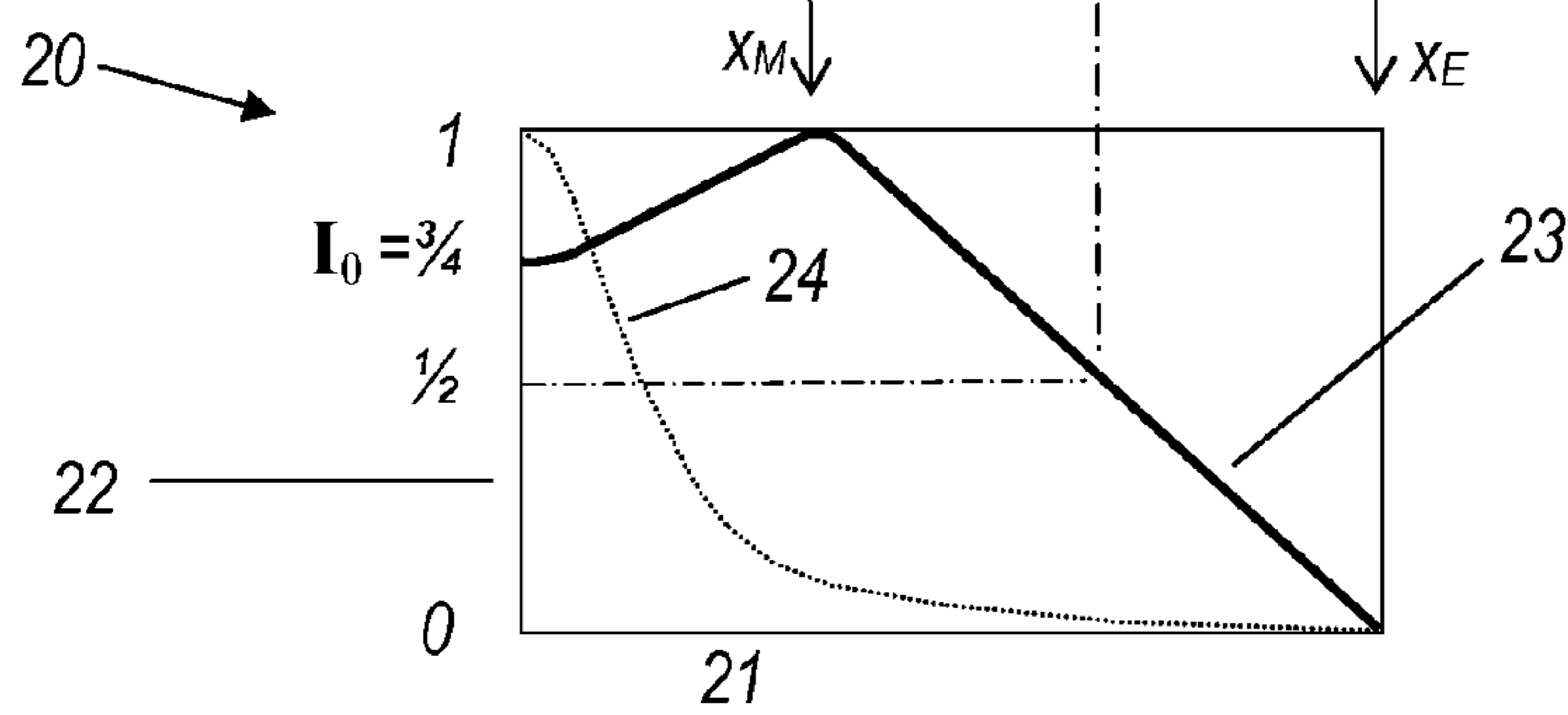
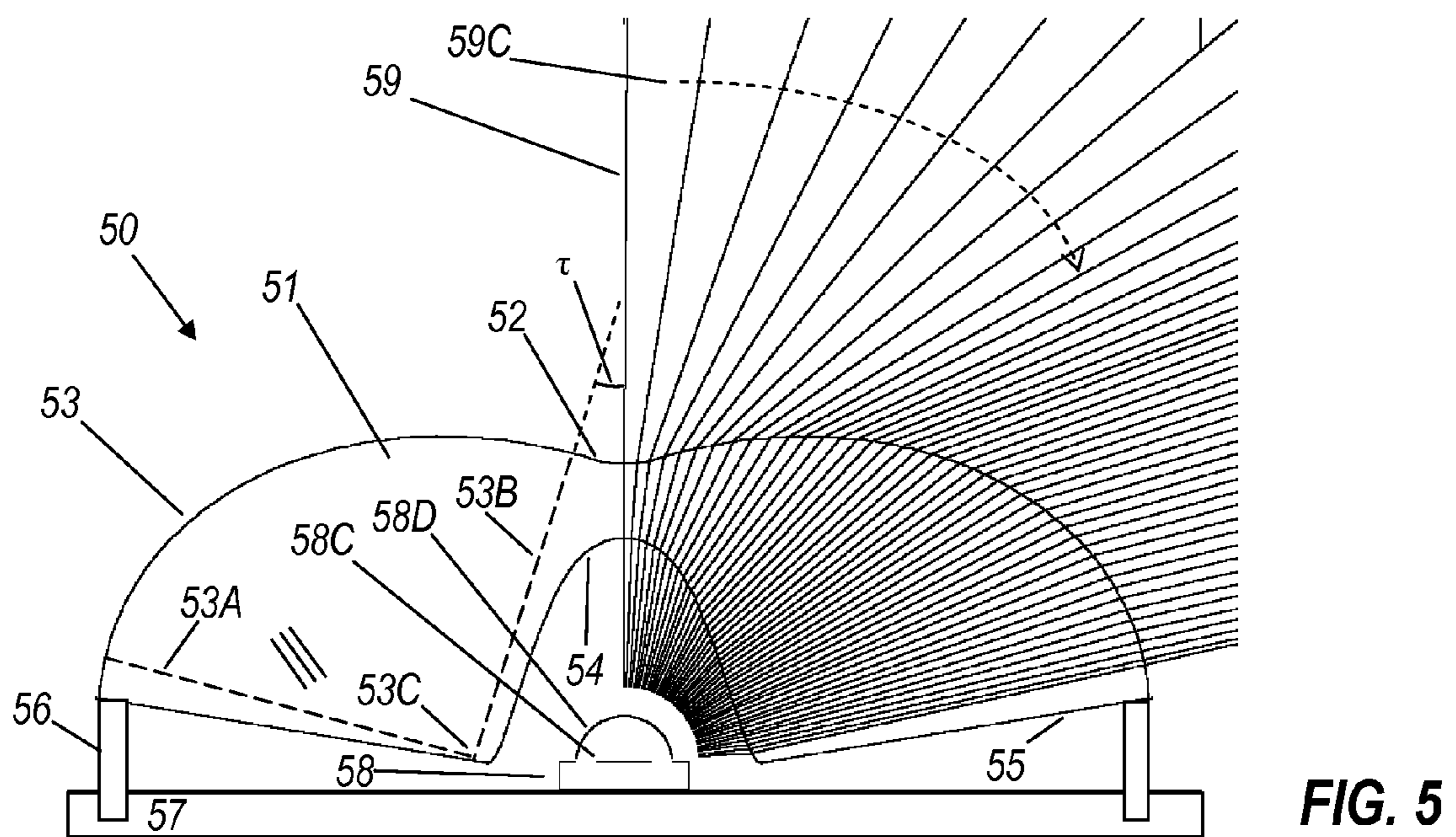
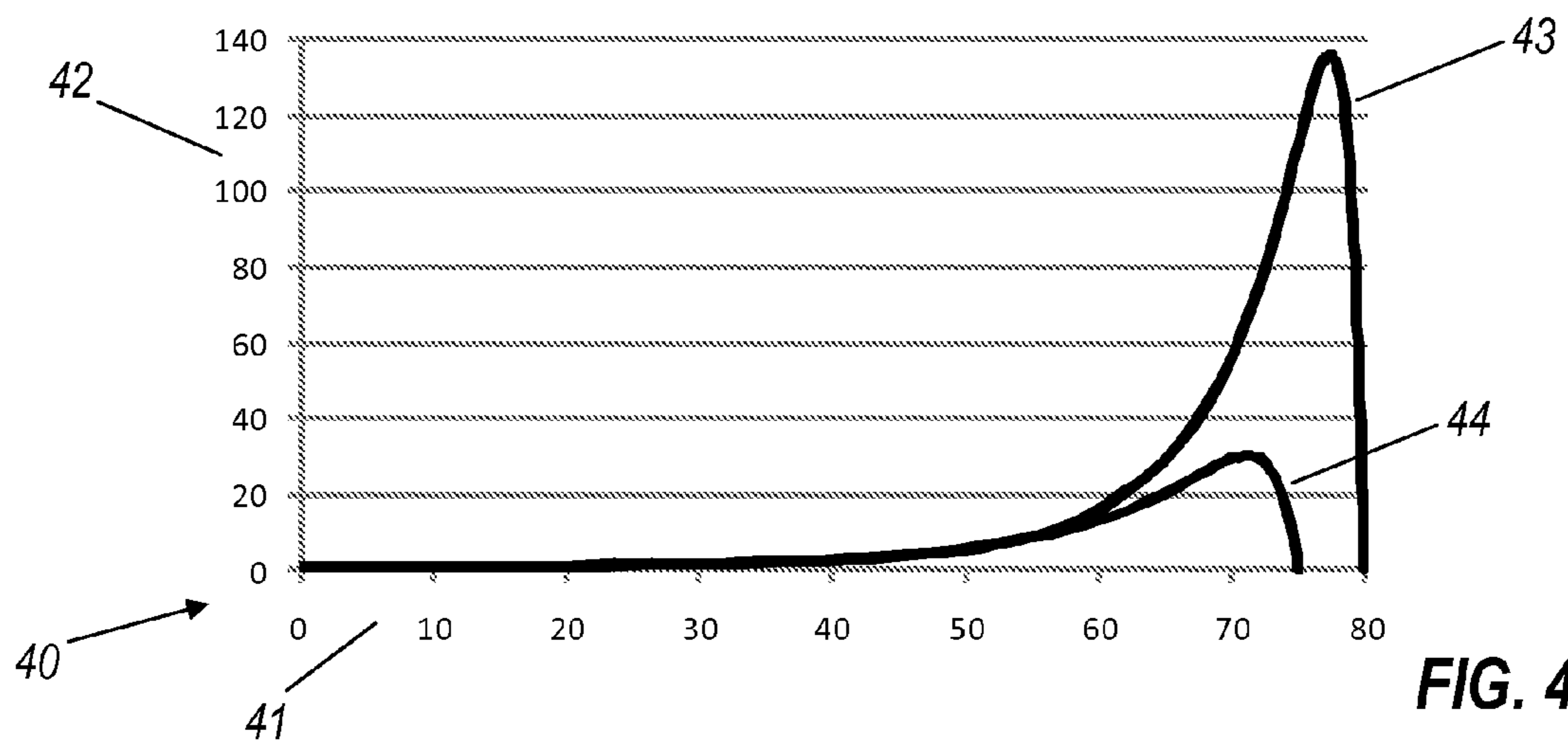
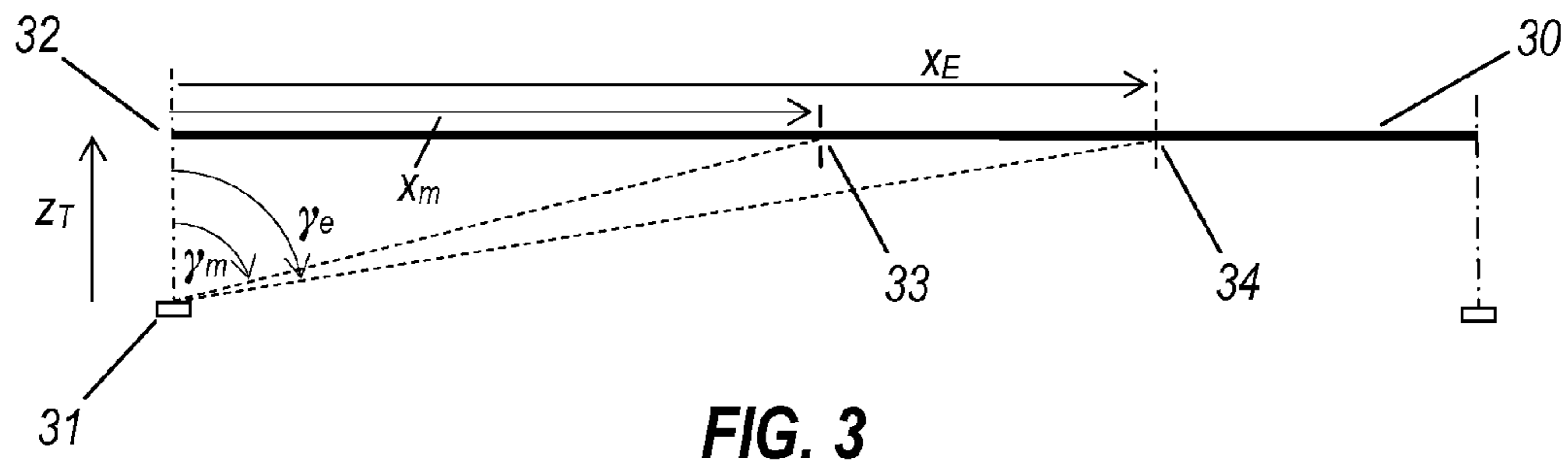


FIG 2



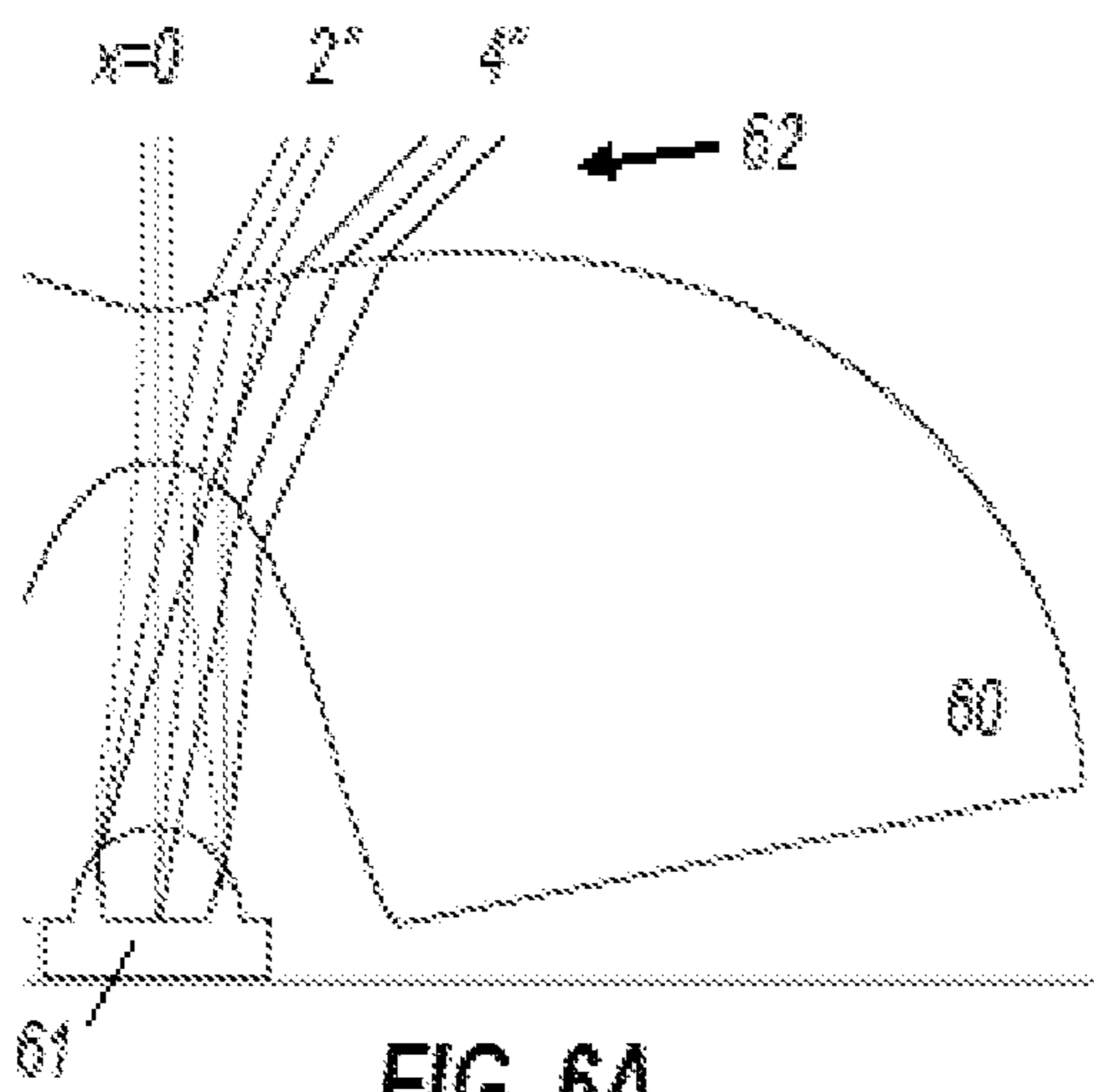


FIG. 6A

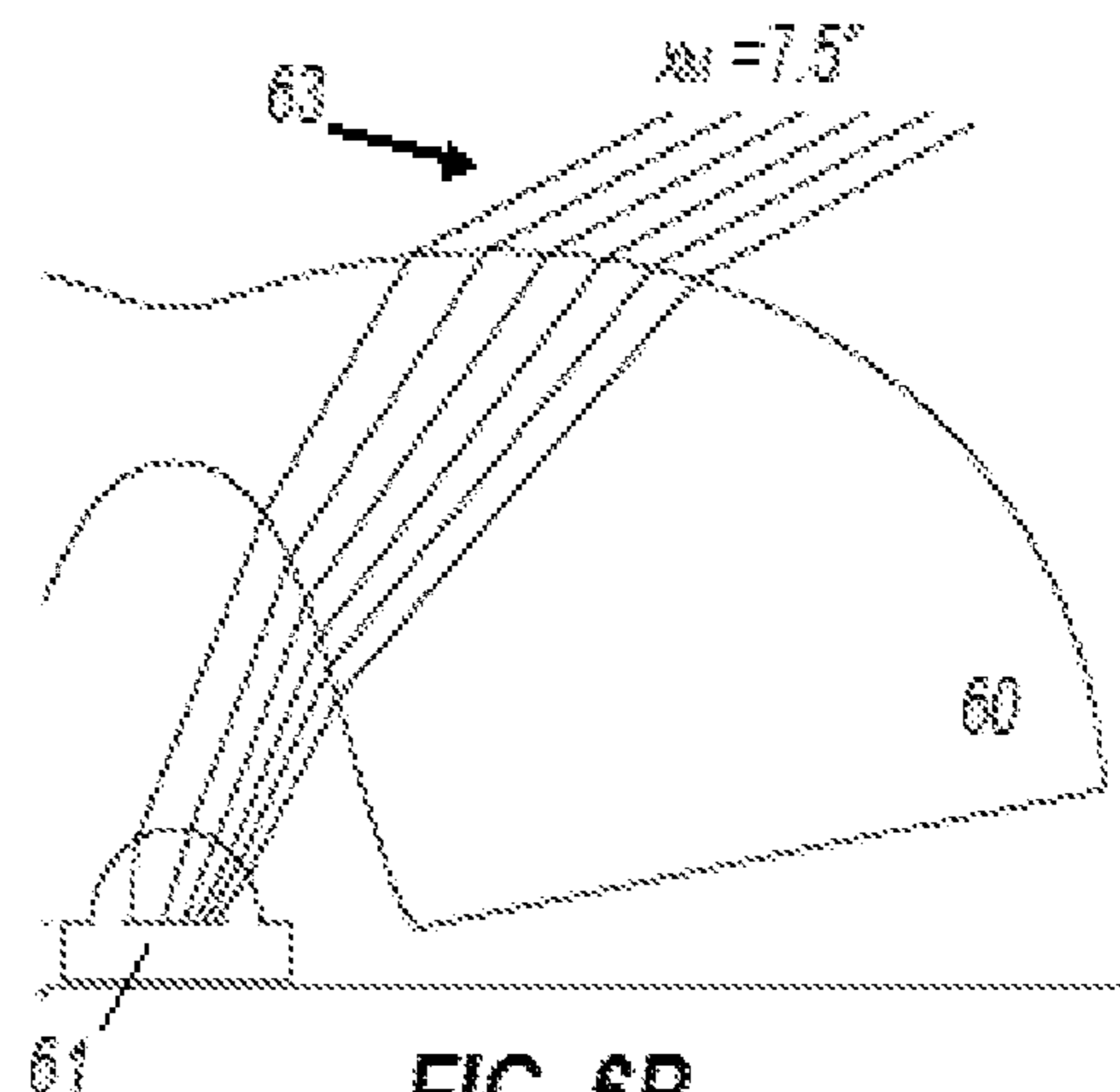


FIG. 6B

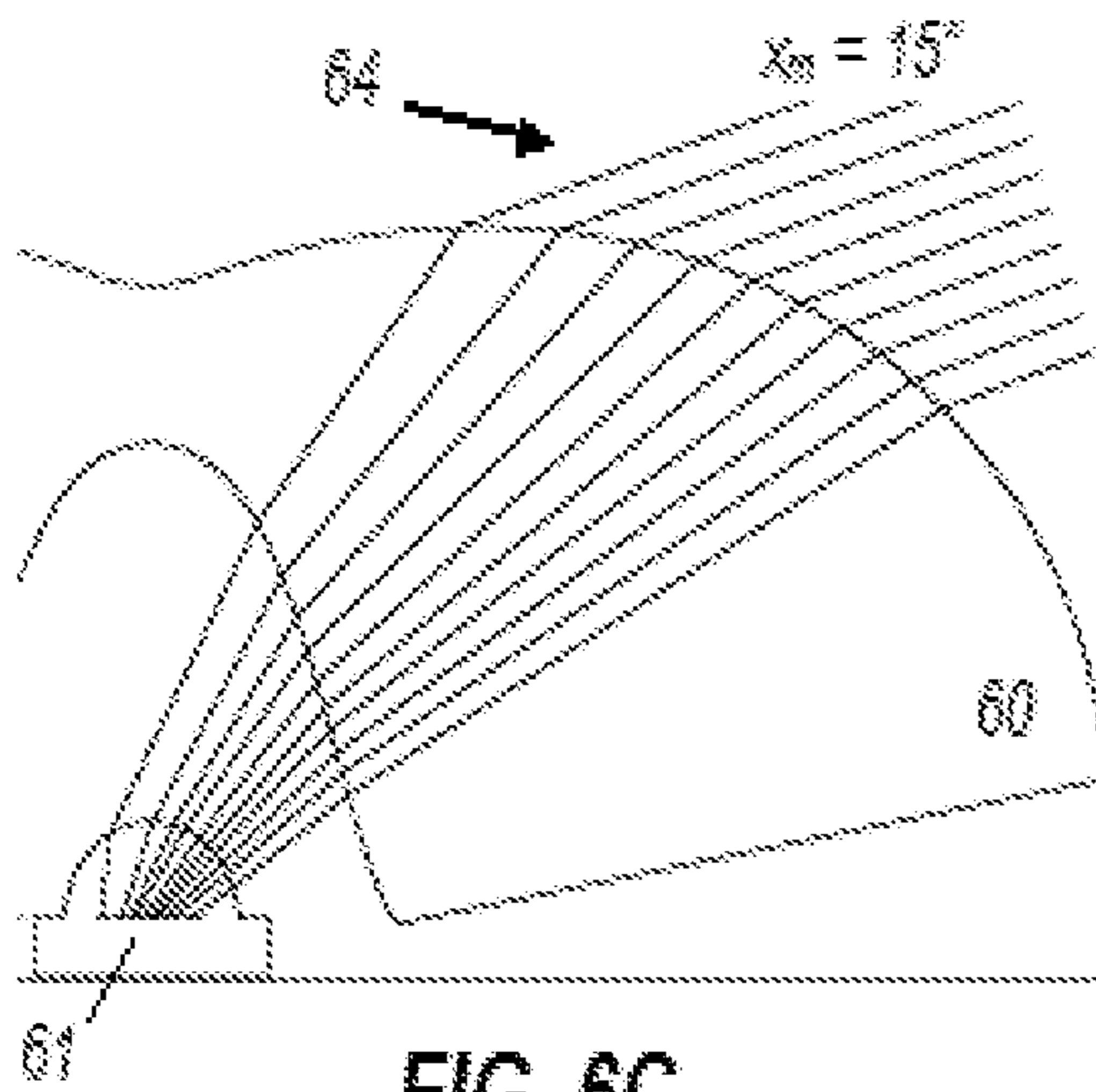


FIG. 6C

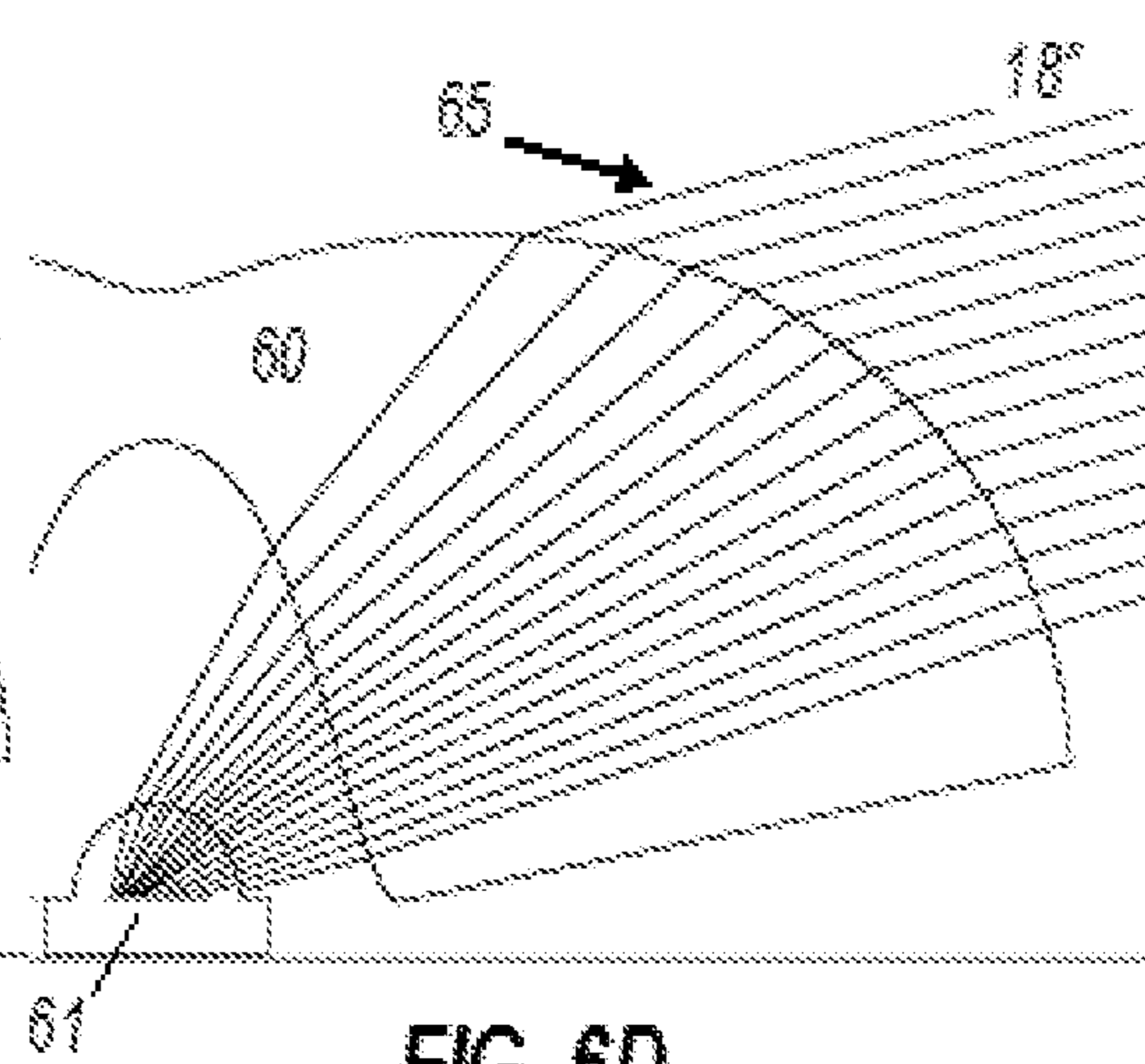


FIG. 6D

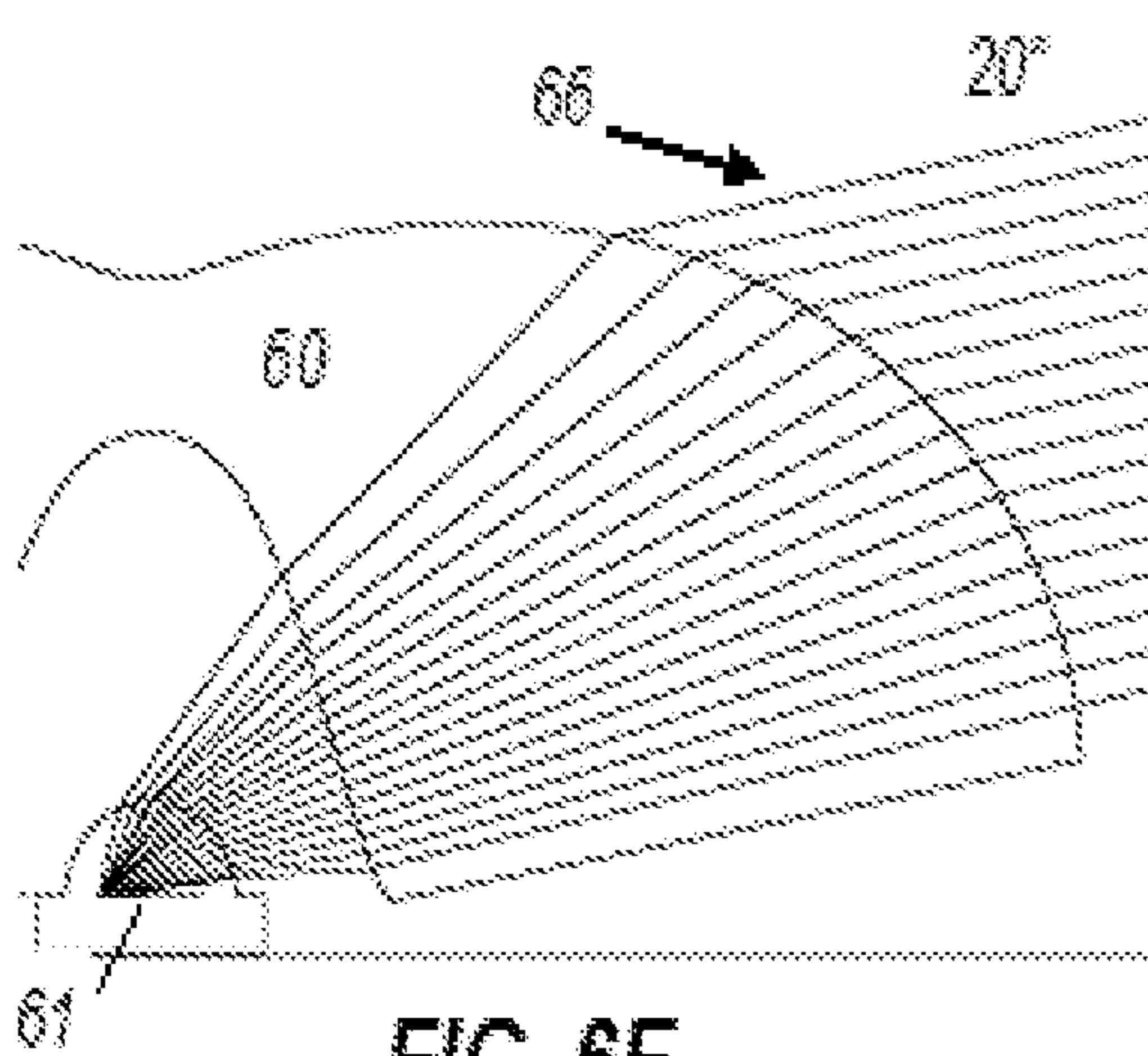


FIG. 6E

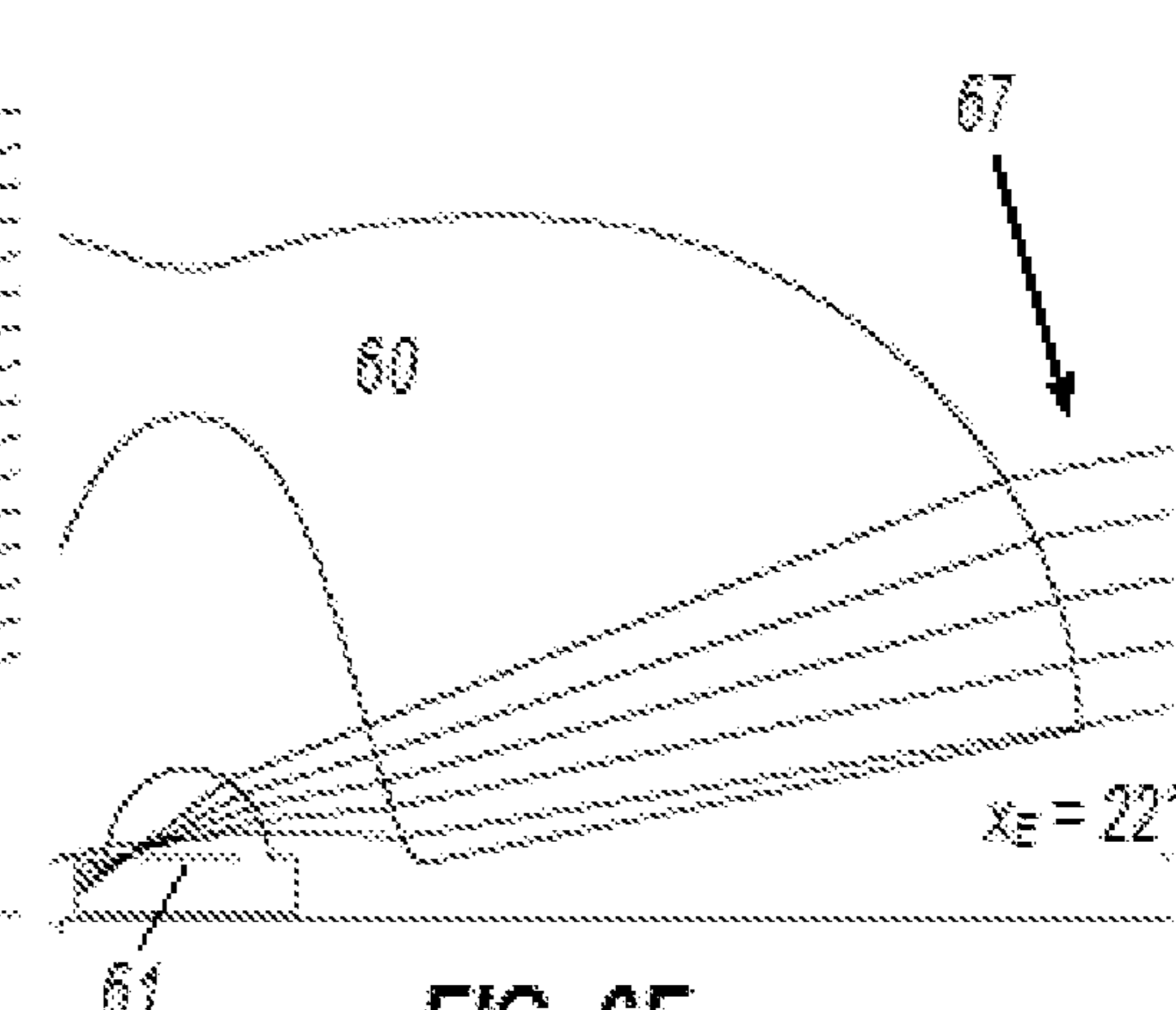


FIG. 6F

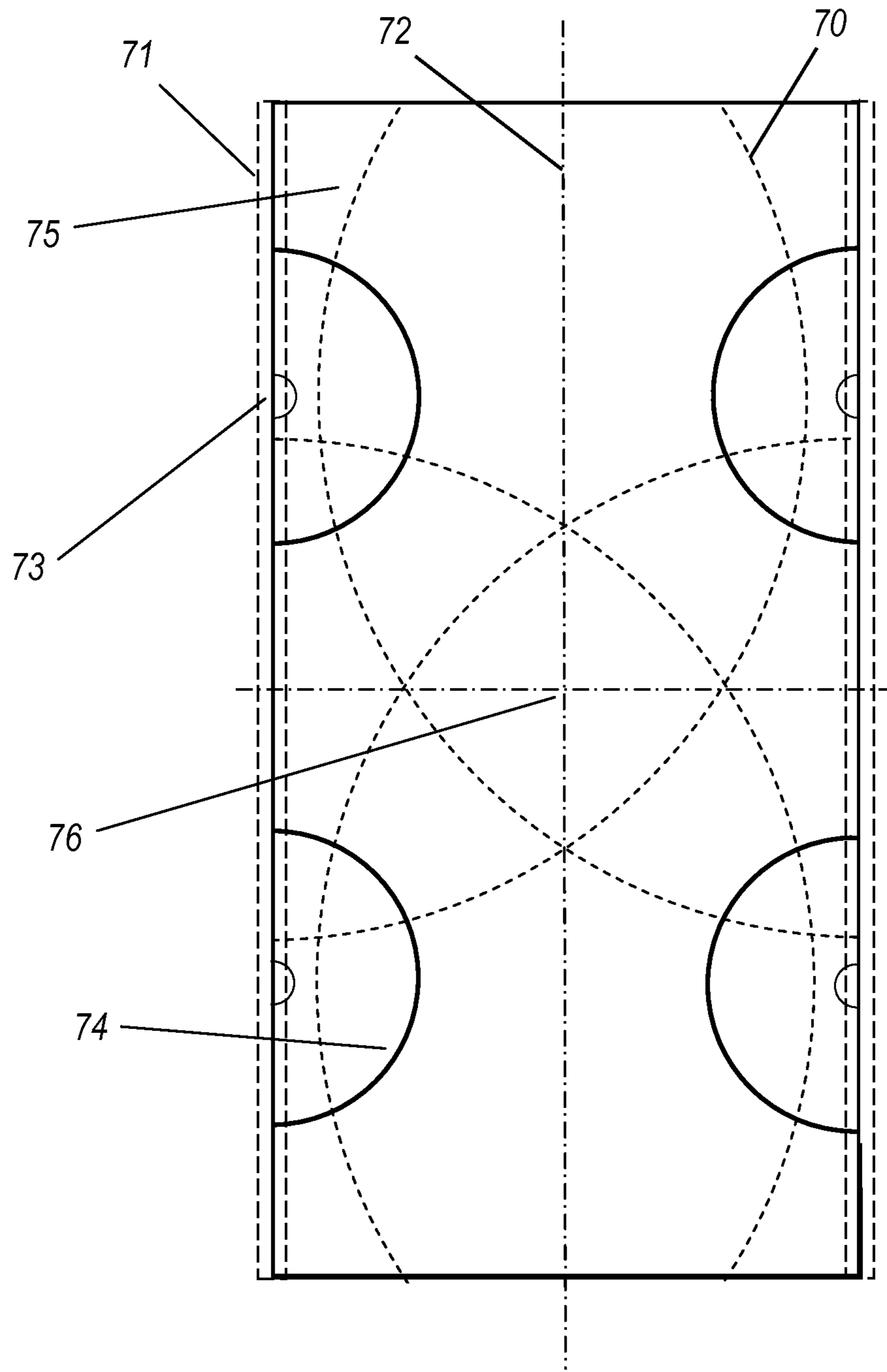
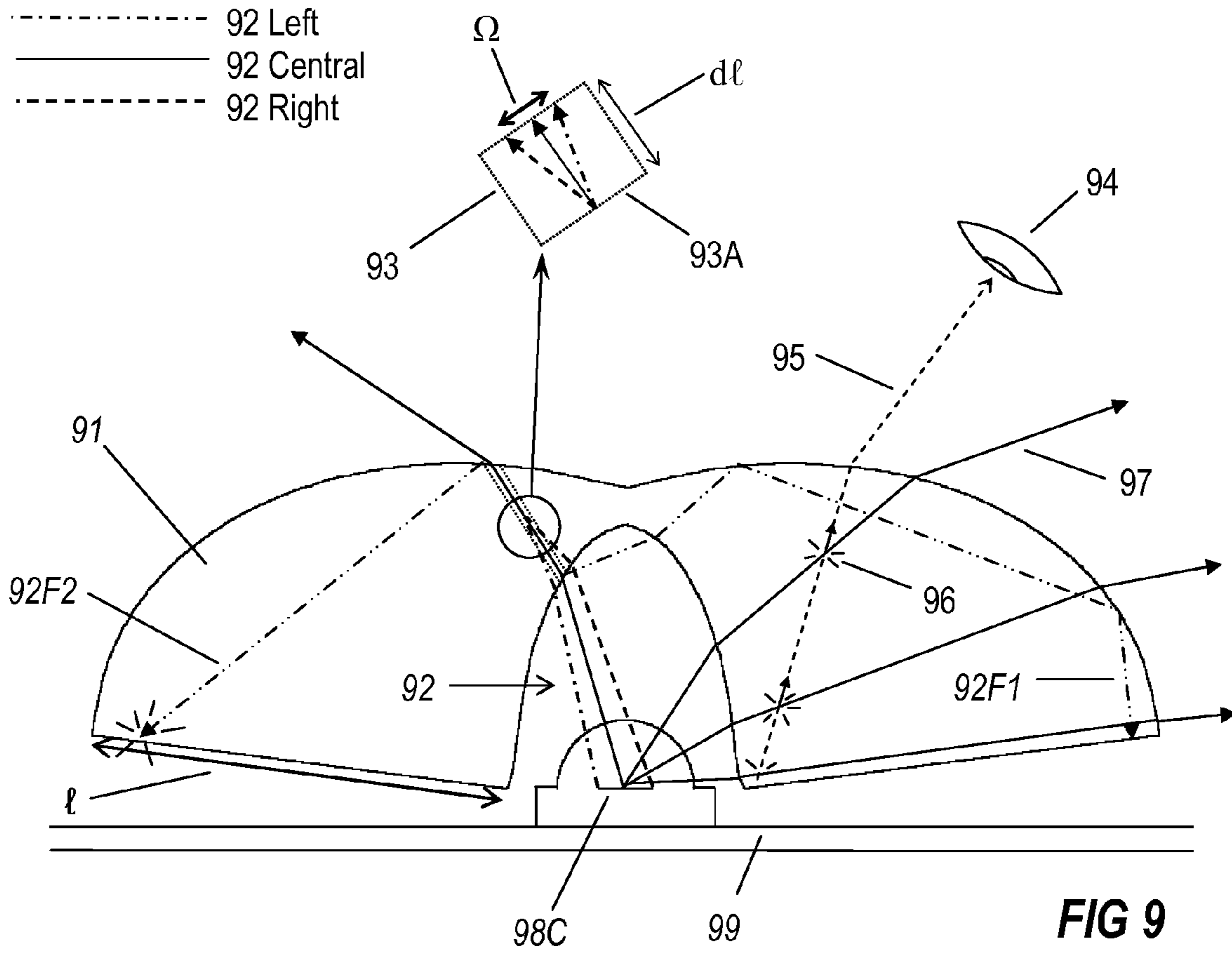
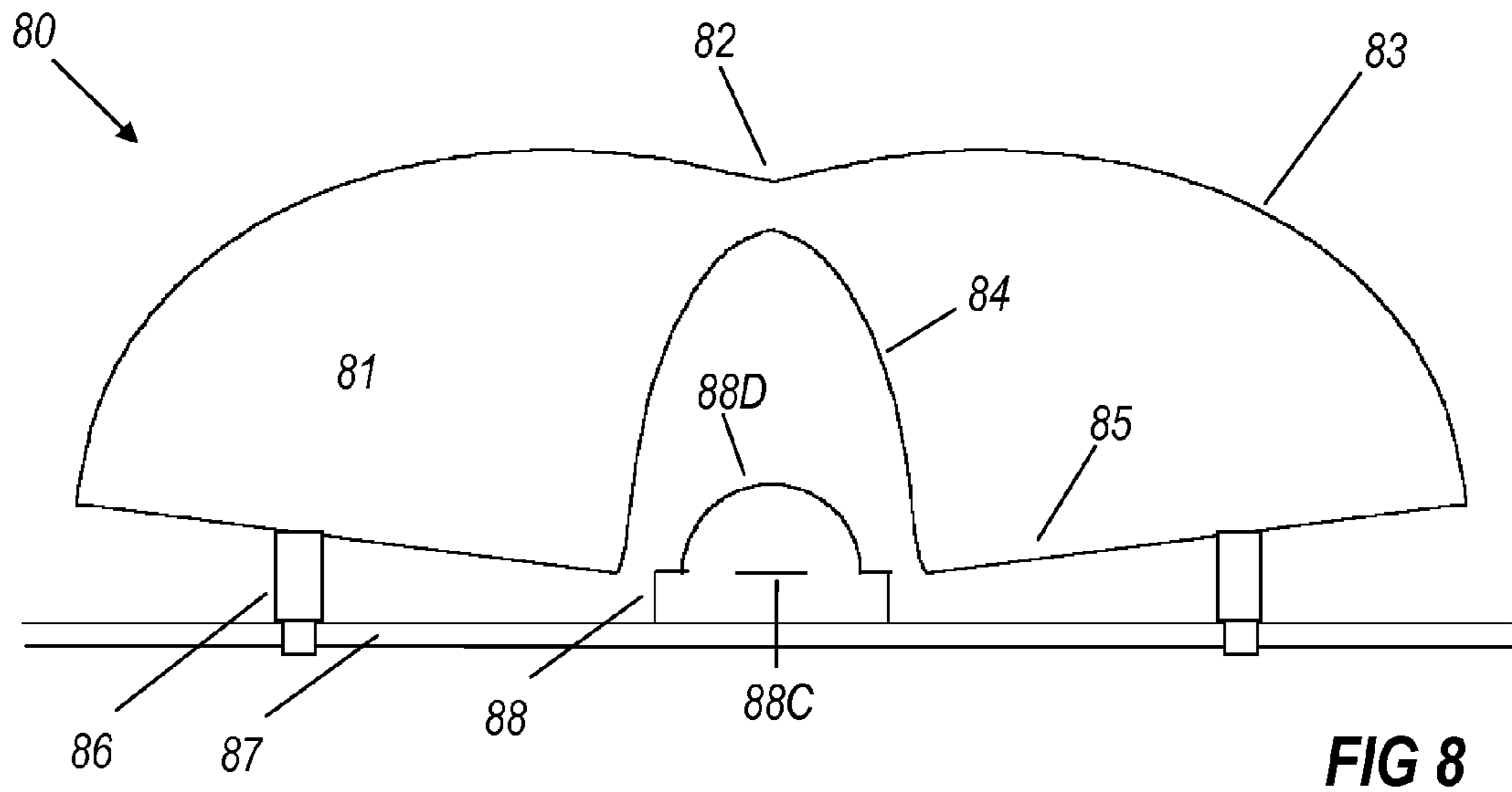
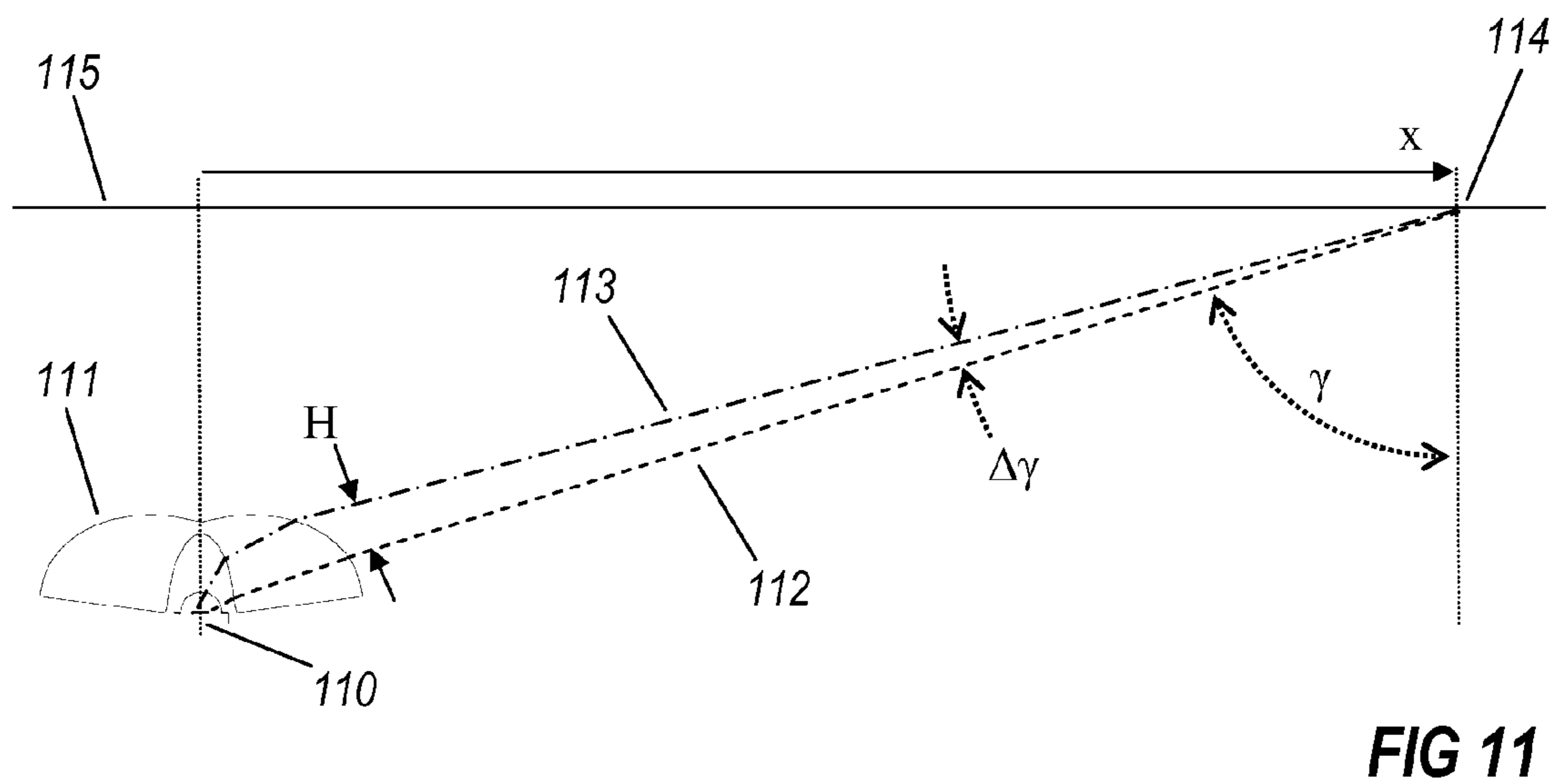
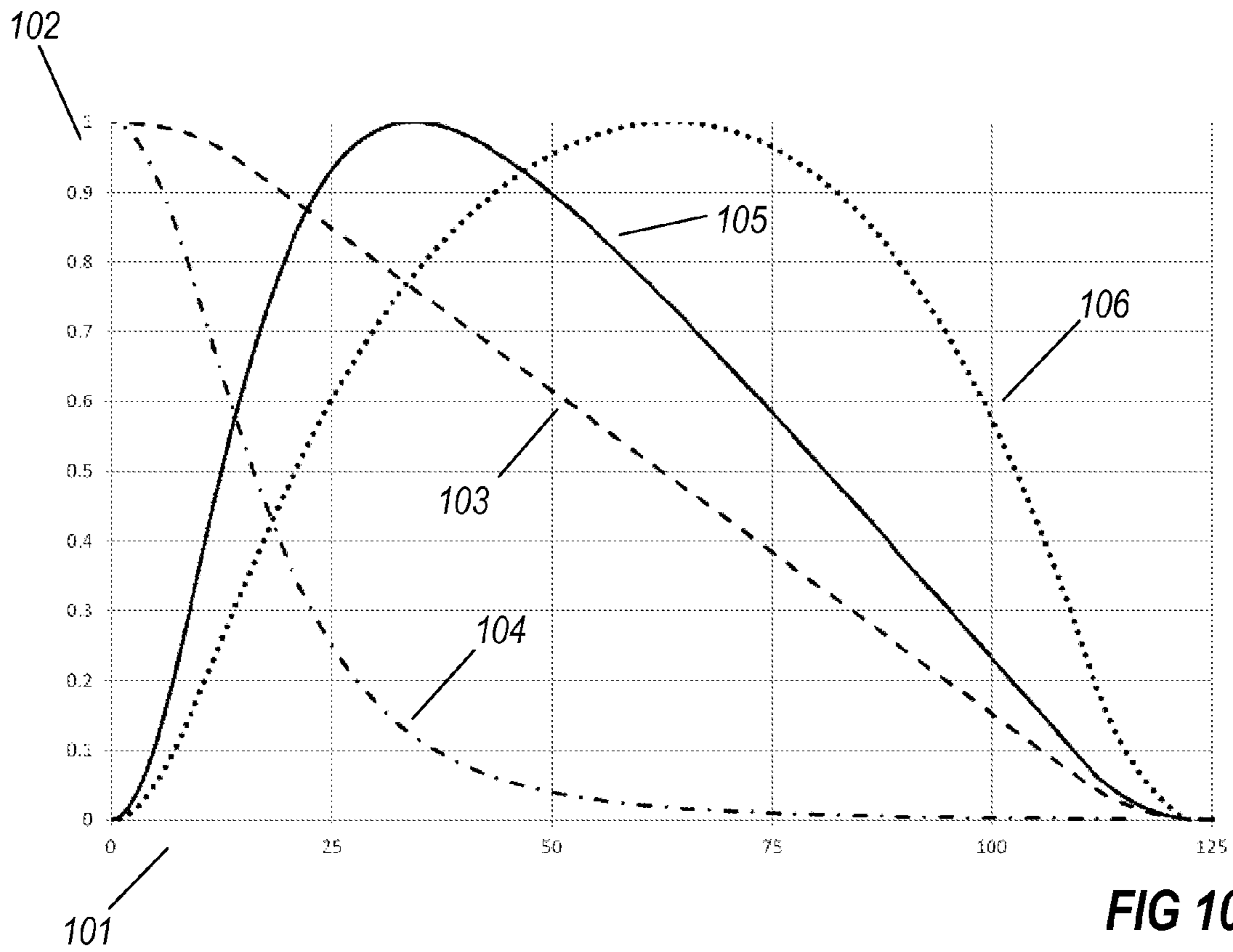


FIG 7





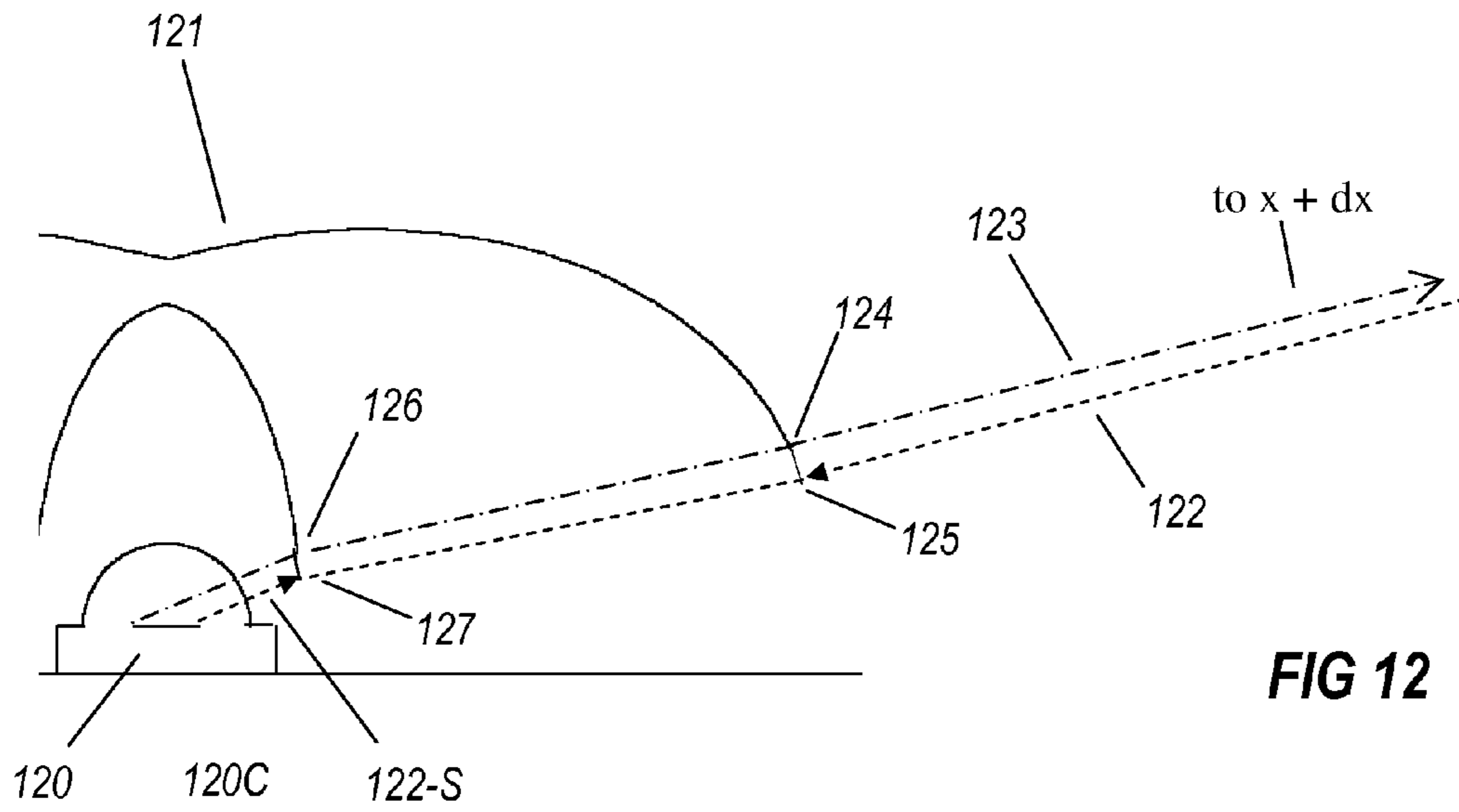


FIG 12

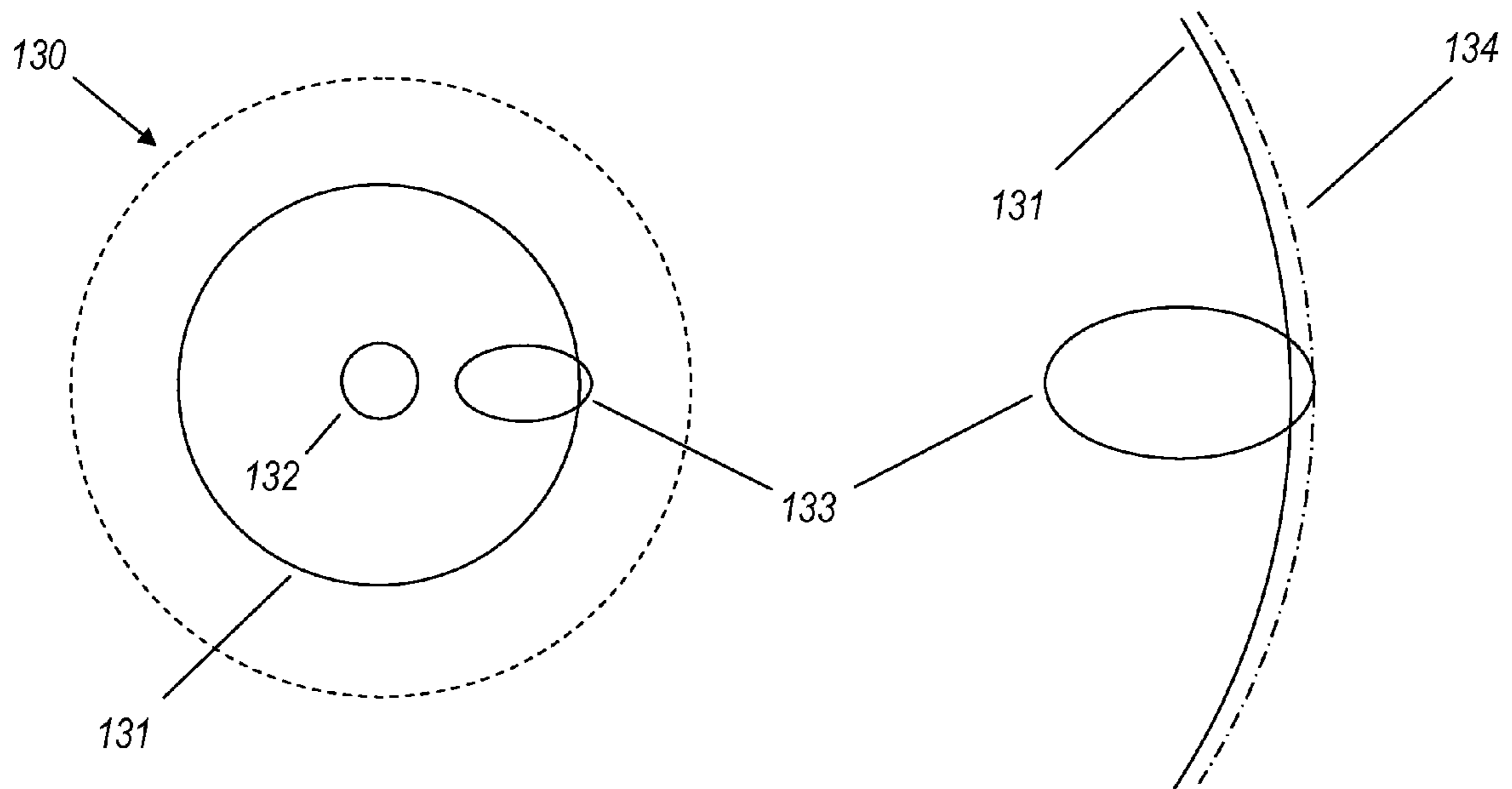


FIG 13A

FIG 13B

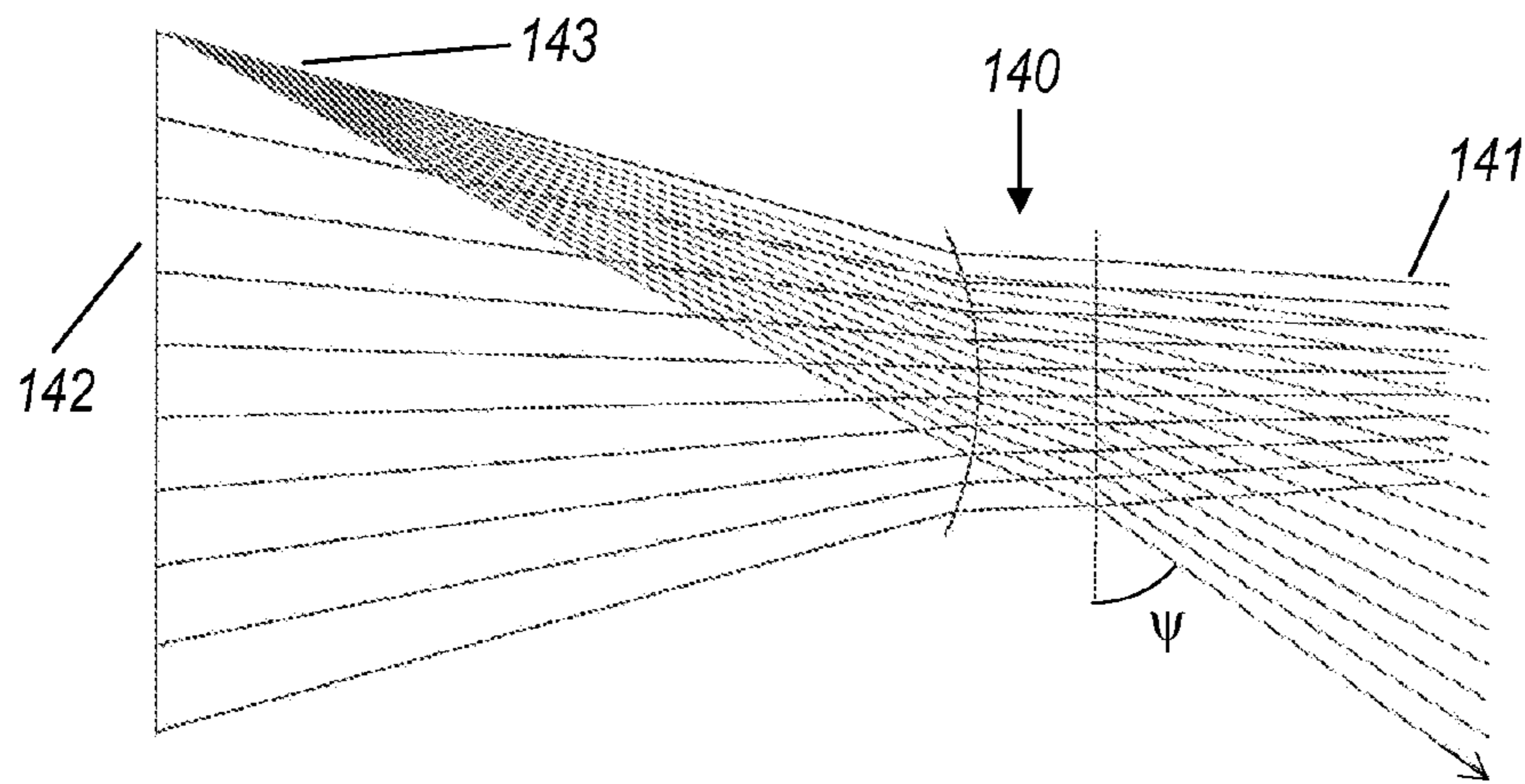


FIG 14

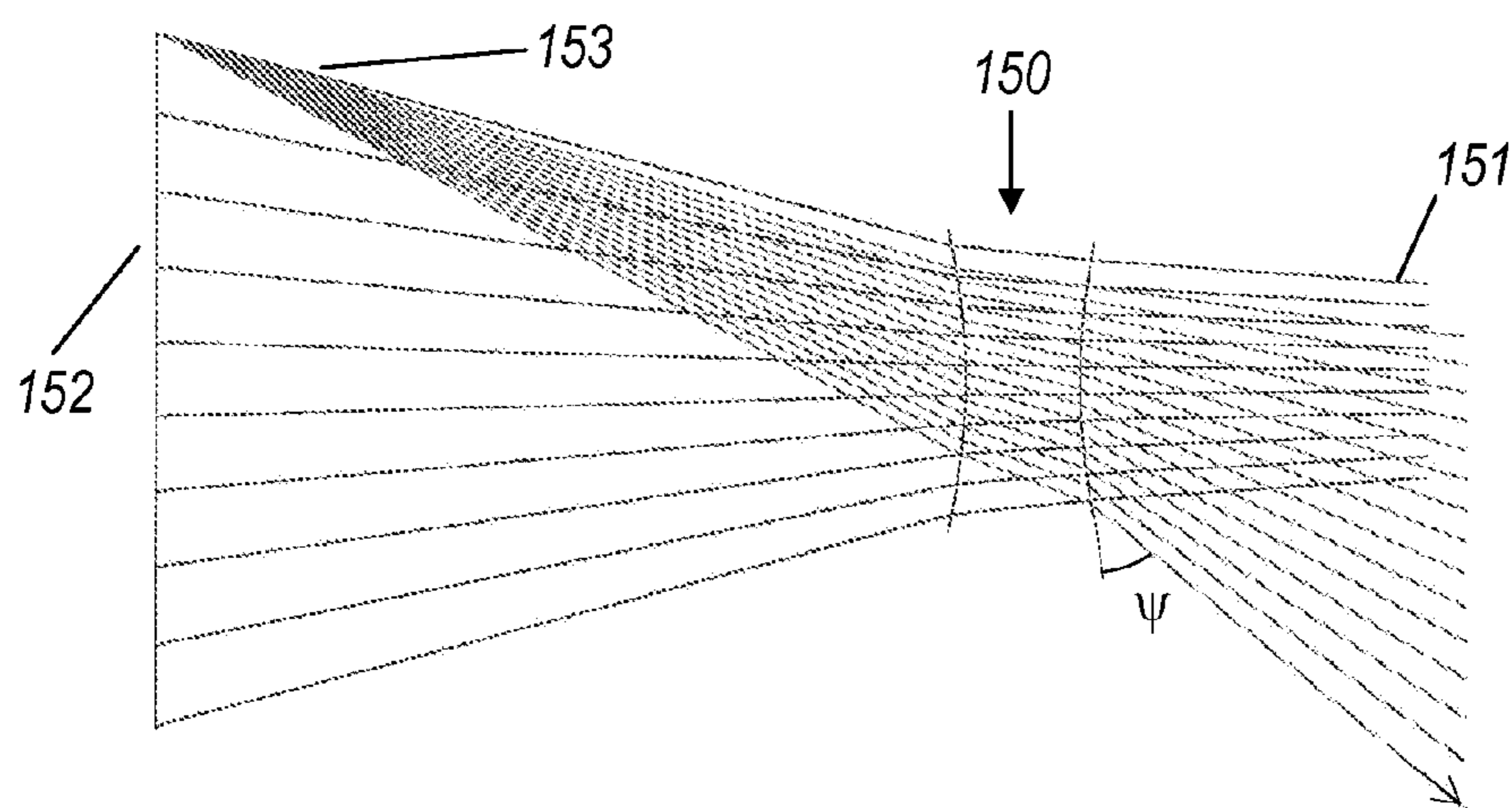


FIG 15

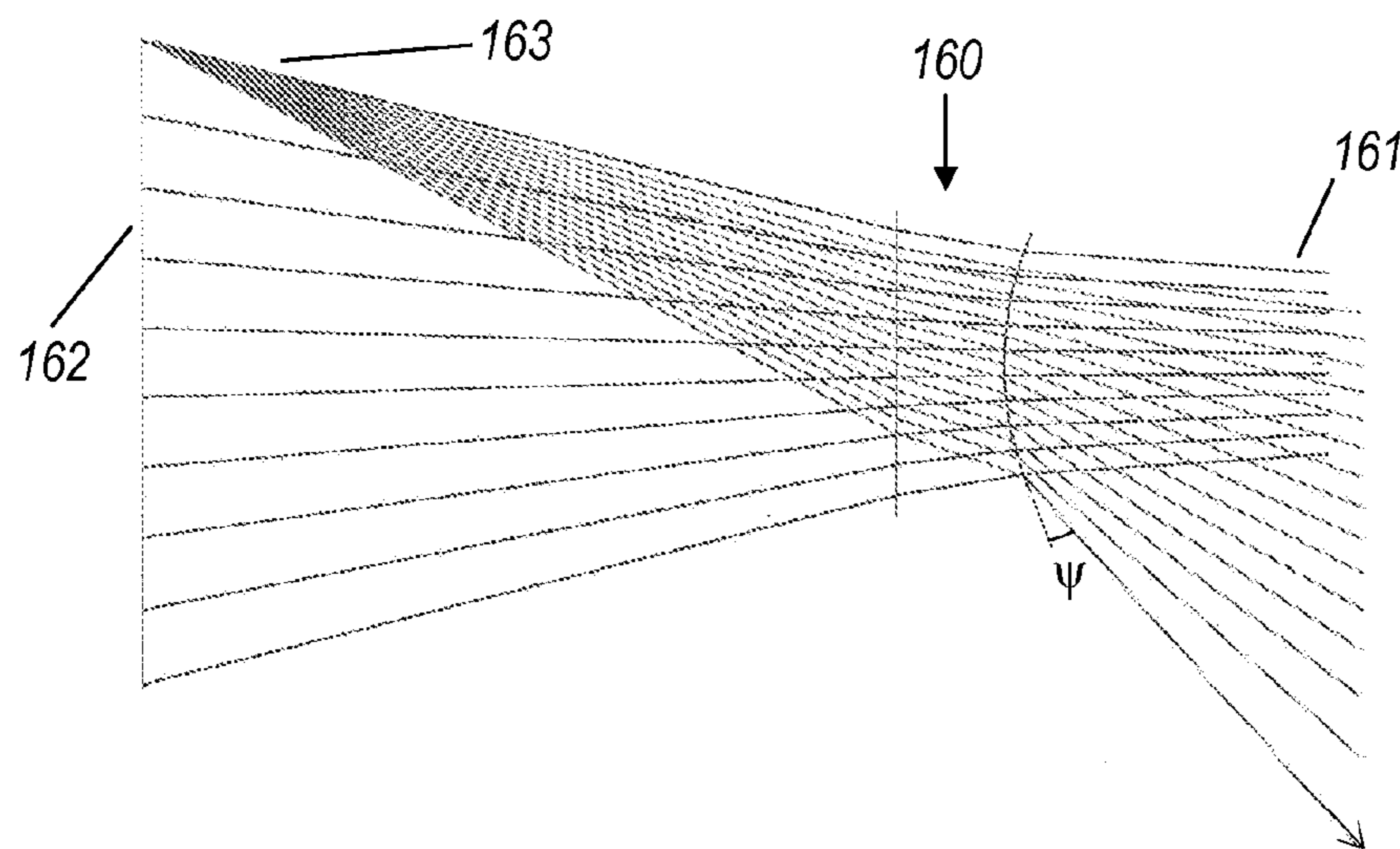


FIG 16

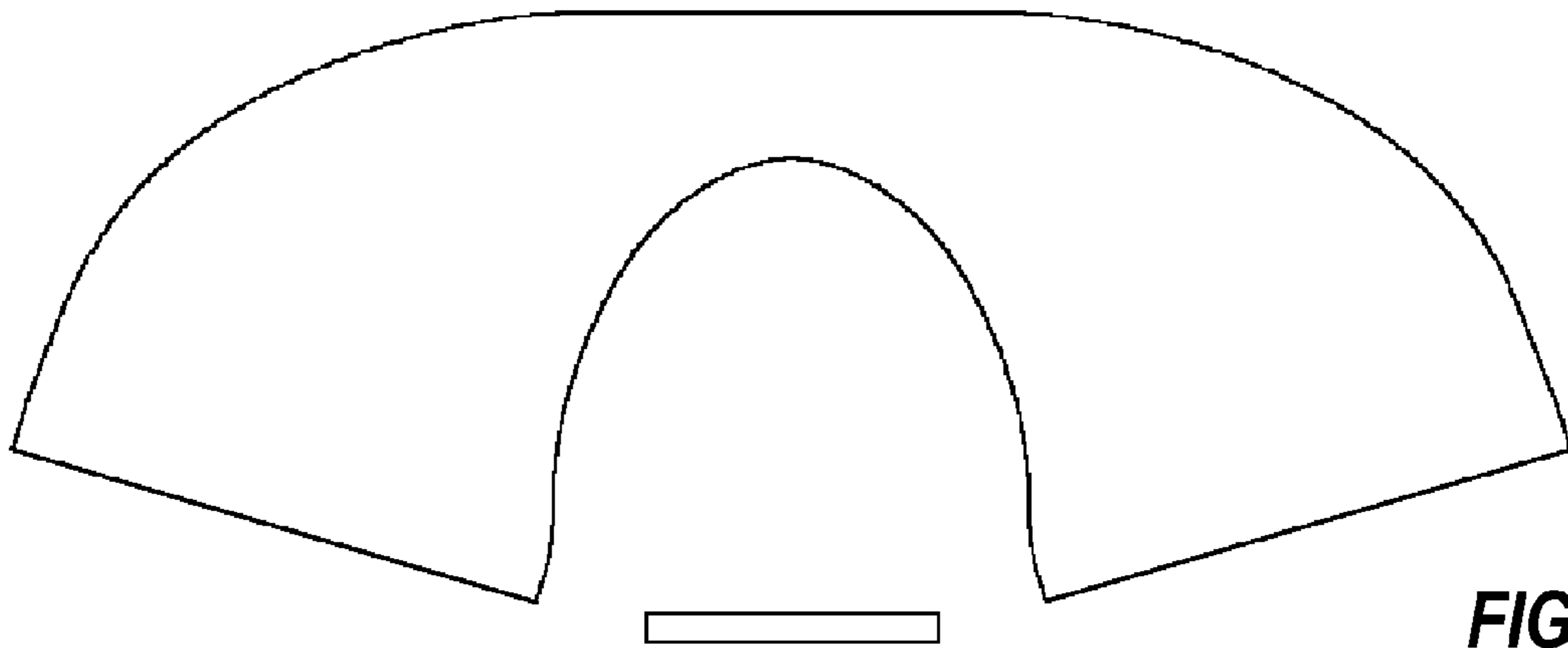


FIG 17

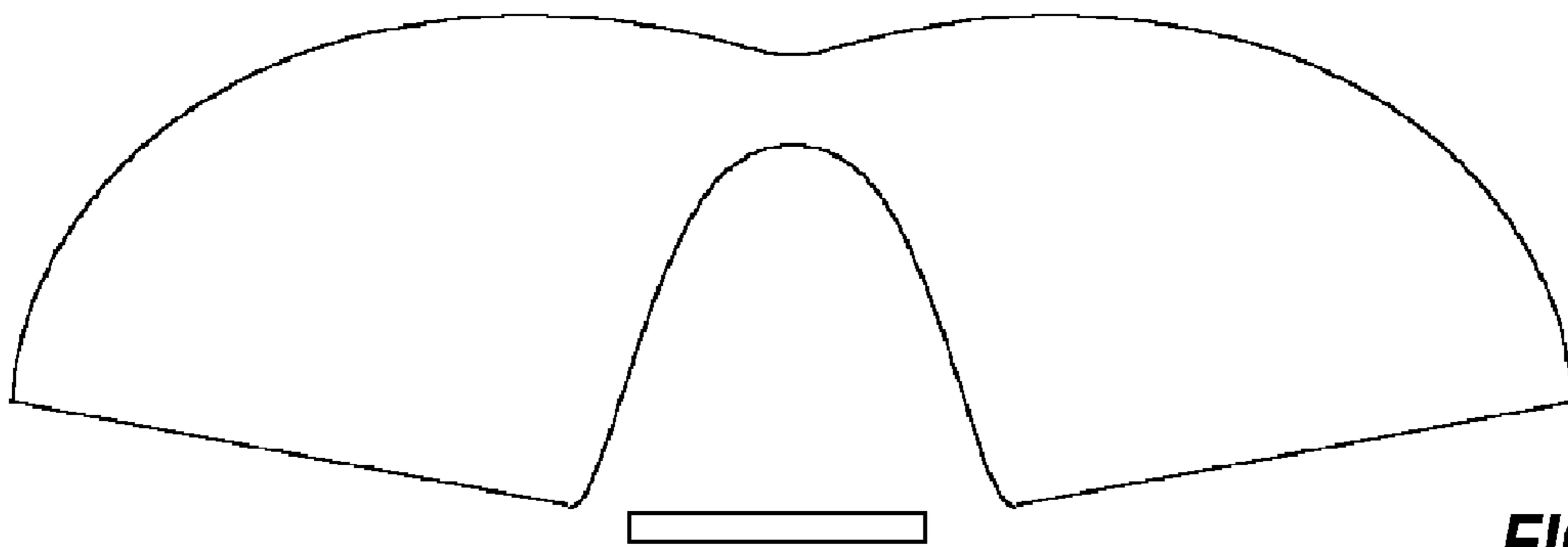


FIG 18

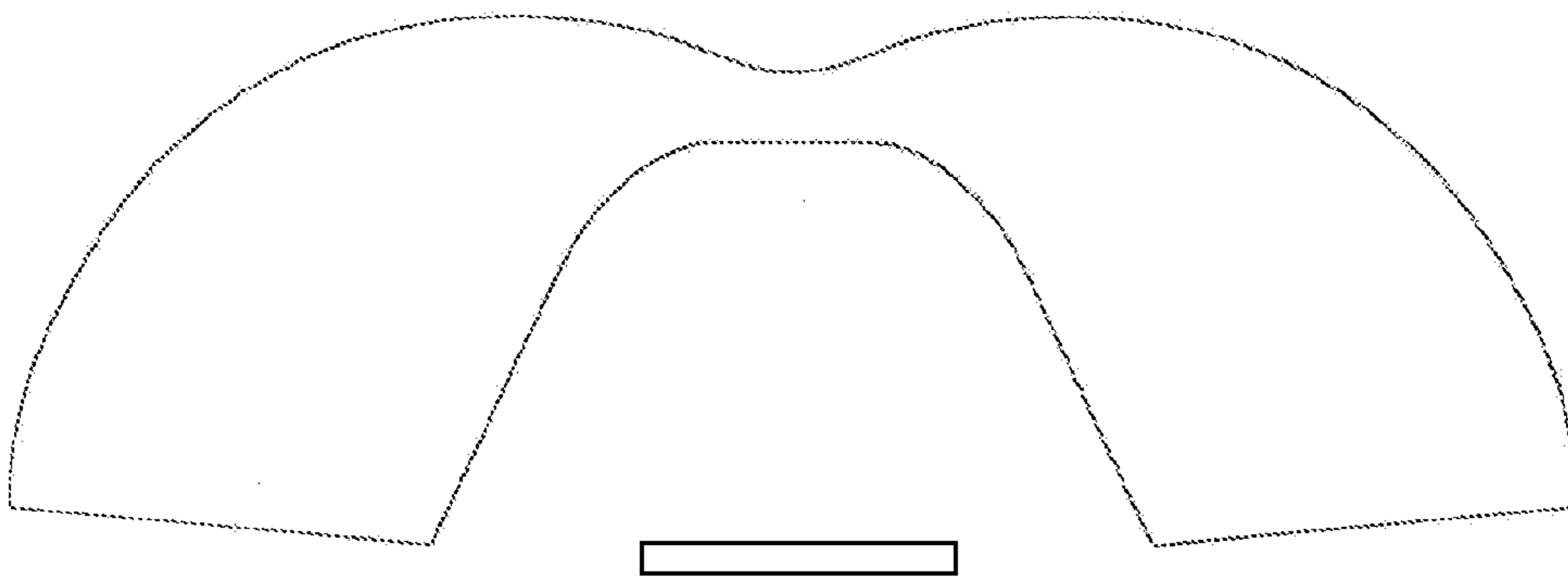


FIG 19

FIG. 20

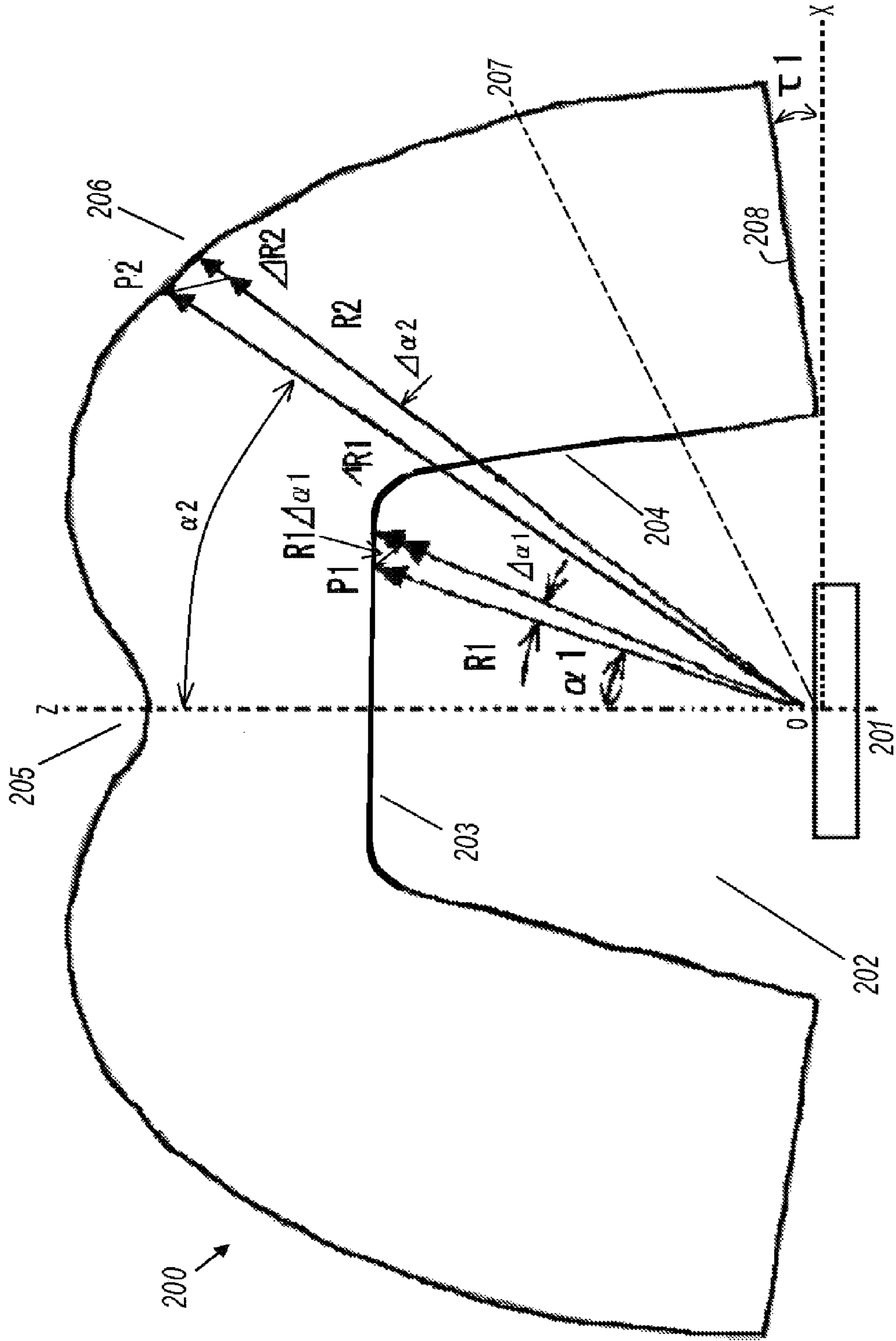
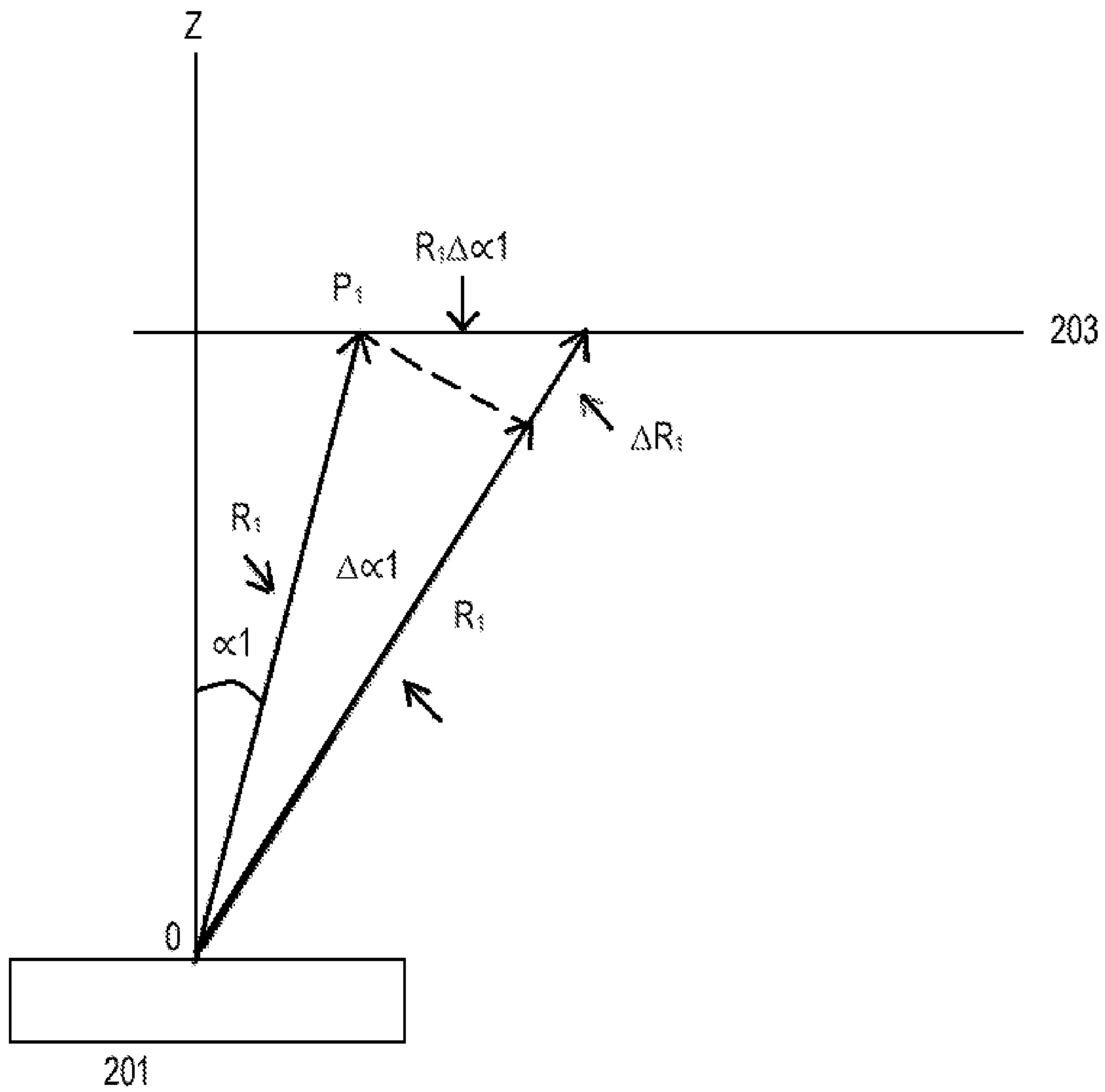


FIG. 21



ILLUMINATION LENS FOR SHORT-THROW LIGHTING

CROSS-REFERENCE TO RELATED APPLICATION

This application claims the benefit of U.S. Provisional Patent Application No. 61/606,710, filed on Mar. 5, 2012, which is hereby incorporated by reference for all purposes as if fully set forth herein.

BACKGROUND

1. Field

Exemplary embodiments of the present invention relate to wide-angle illumination patterns suitable for short-throw lighting, and a general design method for generating their surface profiles.

2. Discussion of the Background

Light emitting diodes (LEDs) may generate light in zones so small (a few mm across) that it is a challenge to spread their flux uniformly over a large target zone, especially one that is much wider than its distance from the LED. So-called short-throw lighting, of close targets, is the polar opposite of spot lighting, which aims at distant targets. Just as LEDs by themselves may not produce a spotlight beam, and so need collimating lenses, they may be equally unsuitable for wide-angle illumination as well, and so need illumination lenses to do the job.

A prime example of short throw lighting is the optical lens for the back light unit (BLU) for a direct-view liquid crystal display (LCD) TVs. Here, the overall thickness of the BLU is usually 26 mm or less and the inter-distance between LEDs is about 200 mm. LCD backlighting may consist of fluorescent tubes arrayed around the edge of a transparent waveguide, that inject their light into the waveguide, which performs the actual backlighting by uniform ejection. While fluorescent tubes may be on the backlight perimeter due to their thickness, light-emitting diodes are so much smaller that they can be placed directly behind the LCD display, (so called "direct-view backlight"), but their punctate nature makes uniformity more difficult, prompting a wide range of prior art over the last twenty years. Not all of this art, however, was suitable for ultra-thin displays.

Another striking application with nearly as restrictive an aspect ratio is that of reach-in refrigerator cabinets. Commercial refrigerator cabinets for retail trade commonly have glass doors with lighting means installed behind the door hinging posts, which in the trade are called mullions. Until recent times, tubular fluorescent lamps have been the only means of shelf lighting, in spite of how cold conditions negatively affect their luminosity and lifetime. Also, fluorescent lamps may produce a very non-uniform lighting pattern on the cabinet shelves. Light-emitting diodes, however, may be favored by cold conditions and are much smaller than fluorescent tubes, which allow for illumination lenses to be employed to provide a much more uniform pattern than fluorescent tubes ever could. Because fluorescent tubes radiate in all directions instead of just upon the shelves, much of their light is wasted. With the proper illumination lenses, however, LEDs can be much more efficient, allowing lower power levels than fluorescent tubes, in spite of the latter's good efficacy.

LED illumination lenses can be classified into three groups, according to how many LEDs are used:

(1) Extruded linear lenses with a line of small closely spaced LEDs, such as described in U.S. Pat. Nos. 7,273,299 and 7,731,395.

(2) Circularly symmetric illumination lenses.

(3) Free-form illumination lenses with rectangular patterns, such as described in U.S. Pat. No. 7,674,019.

The first two approaches may require many LEDs in order to achieve reasonable uniformity, but recent trends in LEDs have produced such high luminosity that fewer LEDs may be needed, allowing significant power savings. This is the advantage of the last approach, but free-form lenses generating rectangular patterns have proved difficult to produce, via injection molding, with sufficient figural accuracy for their overlaps to be caustic-free (caustics are conspicuous small regions of elevated illuminance).

The prior art is even more challenged, moreover, when fewer LEDs are needed due to ongoing year-over-year improvements in LED flux output. After all, backlight thickness is actually relative to the inter-LED spacing, not to the overall width of the entire backlight. For example, in a 1" thick LCD backlight with 4" spacing between LEDs, the lens task is proportionally similar to the abovementioned refrigerator cabinet. Because of the smaller size of an LCD as compared to a 2.5 by 5 foot refrigerator door, lower-power LEDs with smaller emission area may be used, typically a Top-LED configuration with no dome-like silicone lens.

Non-specific design methods for addressing this problem may be found in U.S. Pat. Appl. No. 2006/0138437 and U.S. Pat. Nos. 7,348,723; 7,445,370; 7,621,657; and 7,798,679. U.S. Pat. No. 7,798,679, however, contains only generically vague descriptions of that lens profile, and has no specific method of distinguishing the vast number of significantly different shapes fitting its general description.

Experience has shown that illumination lenses may be unforgiving of small shape errors, such as result from unskilled injection molding or subtle design flaws. Very small changes in local slope of a lens can result in highly visible illumination artifacts sufficient to ruin an attempt at a product. Therefore such generic descriptions, such those in U.S. Pat. No. 7,798,679, may be insufficient for practical use, because even the most erroneous and ill-performing lens fulfills them just as well as an accurate, high-performing lens. Further, U.S. Pat. No. 7,798,679 never provides the specific, distinguishing shape-specifications whose precise details are so necessary for modern optical manufacturing.

SUMMARY OF THE INVENTION

Disclosed herein are embodiments that generate wide-angle illumination patterns suitable for short-throw lighting. Also disclosed is a general design method for generating their surface profiles, one based on nonimaging optics, specifically a new branch thereof, photometric nonimaging optics. This field applies the foundational nonimaging-optics idea of etendue in a new way, to analyze illumination patterns and classify them according to the difficulty of generating them, with difficulty defined as the minimum size lens required for a given size of the light source, in this case, the LED.

It is the first objective of this invention to disclose numerically-specific lens configurations that in arrays will provide uniform illumination for a close planar target, especially in retail refrigeration displays and in thinnest-possible direct-view LCD backlights.

It is the second objective of this invention to provide compensation for the illumination-pattern distortions caused by volume scattering and scattering due to Fresnel reflections, which together act as an additional, indiscriminating secondary light source.

It is to be understood that both the foregoing general description and the following detailed description are exem-

plary and explanatory and are intended to provide further explanation of the invention as claimed.

BRIEF DESCRIPTION OF THE DRAWINGS

The accompanying drawings, which are included to provide a further understanding of the invention and are incorporated in and constitute a part of this specification, illustrate embodiments of the invention, and together with the description serve to explain the principles of the invention.

FIG. 1 shows a rectangular door illuminated by circular patterns.

FIG. 2 shows a graph of an individual illumination pattern.

FIG. 3 shows an end view of the door of FIG. 1, with slant angles.

FIG. 4 shows a graph of required source magnification.

FIG. 5 shows a cross-section of an illumination lens and LED.

FIG. 6A-6F show source-image rays from across the target.

FIG. 7 shows how a rectangular door is illuminated by only 4 LEDs.

FIG. 8 shows a cross-section of a further illumination lens and LED.

FIG. 9 illustrates a mathematical description of volume scattering.

FIG. 10 is a graph of illumination patterns.

FIG. 11 sets up a 2D source-image method of profile generation.

FIG. 12 shows the method of profile generation according to FIG. 11.

FIGS. 13A and 13B show the 3D source-image method of profile generation.

FIG. 14 shows a plano-convex lens-center, with defining rays.

FIG. 15 shows a concave-concave lens-center, with defining rays.

FIG. 16 shows a concave-plano lens-center, with defining rays.

FIG. 17 shows the complete lens made from the lens-center of FIG. 14.

FIG. 18 shows the complete lens made from the lens-center of FIG. 15.

FIG. 19 shows the complete lens made from the lens-center of FIG. 16.

FIG. 20 shows a short throw illumination lens according to FIG. 19.

FIG. 21 shows a detailed view of features of the short throw illumination lens of FIG. 20.

DETAILED DESCRIPTION OF THE ILLUSTRATED EMBODIMENTS

The invention is described more fully hereinafter with reference to the accompanying drawings, in which embodiments of the invention are shown. This invention may, however, be embodied in many different forms and should not be construed as limited to the embodiments set forth herein. Rather, these embodiments are provided so that this disclosure is thorough, and will fully convey the scope of the invention to those skilled in the art. In the drawings, the size and relative sizes of layers and regions may be exaggerated for clarity. Like reference numerals in the drawings denote like elements.

It will be understood that when an element or layer is referred to as being "on" or "connected to" another element or layer, it can be directly on or directly connected to the other element or layer, or intervening elements or layers may be

present. In contrast, when an element is referred to as being "directly on" or "directly connected to" another element or layer, there are no intervening elements or layers present. It will be understood that for the purposes of this disclosure, "at least one of X, Y, and Z" can be construed as X only, Y only, Z only, or any combination of two or more items X, Y, and Z (e.g., XYZ, XYY, YZ, ZZ).

Spatially relative terms, such as "beneath", "below", "lower", "above", "upper" and the like, may be used herein for ease of description to describe one element or feature's relationship to another element(s) or feature(s) as illustrated in the figures. It will be understood that the spatially relative terms are intended to encompass different orientations of the device in use or operation in addition to the orientation depicted in the figures. For example, if the device in the figures is turned over, elements described as "below" or "beneath" other elements or features would then be oriented "above" the other elements or features. Thus, the exemplary term "below" can encompass both an orientation of above and below. The device may be otherwise oriented (rotated 90 degrees or at other orientations) and the spatially relative descriptors used herein interpreted accordingly.

Commercial refrigerator display cabinets for retail sales have a range of distances from mullion to the front of the shelves, commonly from 3" to 8", with the smaller spacing becoming more prevalent as store owners seek to cram more and more product into their reach-in refrigerator cabinets. Fluorescent tubes have great difficulty with these tighter spacing, leading to an acceleration in the acceptance of LED lighting technology. Even though fluorescent tubes is have efficacy comparable to current LEDs, their large size and omnidirectional emission hamper their efficiency, making it difficult to adequately illuminate the mid-shelf region. Early reach-in refrigerator LED illuminators utilized a large number of low-flux LEDs, but continuing advances in luminosity enable far fewer LEDs to be used to produce the same illuminance. This places a premium on having illumination lenses that when arrayed will sum up to uniformity while also having the smallest possible size relative to the size of the LED.

FIG. 1 shows rectangular outline 10 representing a typical refrigerator door that is 30" wide and 60" high, with other doors, not shown, to either side. Dashed rectangles 11 denote the mullions behind which the shelf lighting is mounted, typically at 3-6" from the front of the illuminated shelves. This is much closer than the distance to the shelf center, denoted by centerline 12. There are twelve illuminators (six on either side), four of which are denoted by small circles 1. Each illuminator produces an illuminated circle with its peak on a ring denoted by solid circles 2 and its edge on dotted circles 3. Here the circles 2 have radius of a quarter of the shelf width, or halfway to centerline 12. The circles 3, where illuminance has fallen to zero, are sized to meet the circles 2 from the opposite mullion.

FIG. 2 shows graph 20 with abscissa 21 that is horizontally scaled the same as FIG. 1 above it. Ordinate 22 is scaled from 0 to 1, denoting the ideal illuminance $I(x)$, as graphed by curve 23, generated on the shelves by an illuminator under the mullion. This illumination function is relative to the maximum on circle 2, which has radius x_M . It falls off to zero at radius x_E . This gradually falling illuminance is paired with the gradually ascending one of the illuminator on the opposite side of the door, so the two patterns add up to constant illuminance along the line 4 of FIG. 1.

An actual injection-molded plastic lens will exhibit volume scattering within its material, making the lens itself an emitter rather than a transmitter. This volume scattered light

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will be strongest just over the lens. The central dip in the pattern **23**, shown in FIG. **2** to be at the $3/4$ level, compensates for this extra volume-scattered light, so that the total pattern (direct plus scattered) is flat within circle **2**. This effect becomes more pronounced with the larger lenses discussed below.

Another advantage of this type of gradually falling-off pattern is that any point on centerline **12** is lit by several illuminators on each mullion, assuring good uniformity. The dotted curve **24** shows the illumination pattern of an LED alone. It is obviously incapable of adding up to satisfactory illumination, let alone uniform, hence the need for an illumination lens to spread this light out properly.

FIG. **3** shows an end view of shelf-front rectangle **30** identical to that of FIG. **1**. Illuminators are located as shown by small rectangles **31**. The strategy of this invention addresses the difficulty of lighting from so close to the shelf, in this case at a distance of $z_T=4"$. FIG. **3** shows the distances $x_m=15"$ and $x_E=22.5"$, respectively, to centerline **33** and edge-line **34**, at which the pattern of illuminator **31** has reached zero illuminance. These distances correspond to off-axis angles from the normal given by

$$\gamma_m = \tan^{-1}(x_m/z_T) = \tan^{-1}(15/4) = 75^\circ \quad \gamma_E = \tan^{-1}(x_E/z_T) = \tan^{-1}(22.5/4) = 80^\circ$$

These large slant-angles drive the lens design, requiring considerable lateral magnification of the source by the lens. At low slant-angles, in contrast, the lens must demagnify.

This concept of magnification and demagnification can be made more explicit via etendue considerations. The source-etendue is that of a chip of area $A_S=2.1 \text{ mm}^2$, immersed in a dome of refractive index $n=1.45$:

$$E_S = \pi n^2 A_S \sin^2 \theta = 14 \text{ mm}^2$$

Here θ is 90° for a Lambertian source of which an LED is a very good approximation.

An illumination lens basically redistributes this etendue over the target, which is much larger than the chip. In the case of the illumination pattern in FIG. **2**, the target etendue relates to the area A_T of the 45" illumination circle of FIG. **1**, as weighted by the relative illumination function **23** of FIG. **2**. This simple model of an actual illumination pattern has a central dip to illuminance I_0 , a rise to unity at $x=x_M$, and a linear falloff to zero at $x=x_E$. This is mathematically expressed as

$$I(x) = I_0 + x(1-I_0)/x_M \quad x \leq x_M$$

$$I(x) = (x_E - x)/(x_E - x_M) \quad x_M \leq x \leq x_E$$

Then the target etendue is given by an easily solved integral:

$$E_T = \pi \sin^2 \theta_T \int_0^{x_E} 2\pi x I(x) dx =$$

$$2\pi^2 \sin^2 \theta_T \left(\left[\frac{I_0 x^2}{2} + \frac{(1-I_0)x^3}{3x_M} \right]_0^{x_M} + \left[\frac{x_E}{2(x_E - x_M)} x^2 - \frac{x^3}{3(x_E - x_M)} \right]_{x_M}^{x_E} \right)$$

$$E_T = \sin^2 \theta_T 1.47 \text{ square meters}$$

Here θ_T is the half angle of a narrow-angle collimated beam with the same etendue as the source, so that

$$\sin^2 \theta_T \sim 1E-5 \quad \theta_T = \pm 0.18^\circ$$

At the center of the lens this is reduced by $3/4$, to $\pm 0.13^\circ$. This can be contrasted with the angular subtense of the source alone, as seen from directly above it on the shelf, at distance z_T as shown in FIG. **3**:

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$$\tan^2 \theta_S = n^2 A_C / 4z_T^2 \theta_S = \pm 0.61^\circ$$

Thus the central demagnification of the lens needs to be 1:4.5, dictating that the central part of the lens be concave, in order to act as an expander with negative focal length. This can be attained on a continuum of concavity bounded by a flat-topped outer surface with a highly curved inside surface or a flat-topped inner surface with the outer surface highly curved. That of FIG. **5** lies between these extremes.

As shown in FIG. **3**, a high slant angle γ means that to achieve uniform illumination the source image made by the lens must be correspondingly larger than for normal incidence, by a factor of $1/\cos \gamma$. The source itself will be foreshortened by a slant factor of $\cos \gamma$, as well as looking smaller and smaller by being viewed from farther away, by a further factor of $\cos^2 \gamma$. Thus the required lens magnification is

$$M(\gamma) = \frac{1}{4.5 \cos^4 \gamma}$$

Note that magnification rises from $1/4$ on-axis to unity at an off-axis angle given by

$$\gamma(M) = \cos^{-1} \sqrt[4]{\frac{1}{4.5M}}$$

$$\gamma(1) = 47^\circ$$

These angles dilute the illuminance by a cosine-cubed factor, so that the farther out light must be thrown, the more intense must be the lens output. Considering that the LED source has a cosine fall-off of its own, the total source magnification required is the well-known \cos^{-4} factor, amounting to 223 at 75° 1100 at 80° respectively. Here lies the advantage of the fall-off in the illumination pattern of FIG. **2**, since these deleterious factors are reduced accordingly.

FIG. **4** shows graph **40** with abscissa **41** running from 0 to 80° in off-axis angle γ and ordinate **42** showing the source magnification $M(\gamma)$ required for uniform illumination. Unit magnification is defined as a source image the same size as if there were no lens. What this magnification means is that the illumination lens of the present invention must produce an image of the glowing source, as seen from the shelf, that is much bigger than the Lambertian LED source without any lens. For uniformly illuminating a 4" shelf-distance, curve **43** shows that the required magnification peaks at 77.5° , while lower curve **44** is for the much easier case of a 6" shelf distance, peaking at 71° . This required image-size distribution is the rationale for the configuration of the present invention.

FIG. **5** is a cross-section of illuminator **50**, comprising illumination lens **51**, bounded by an upper surface comprising a central spherical dimple with arc **52** as its profile and a surrounding toroid with elliptical arc **53** as its profile, and also bounded by a lower surface comprising a central cavity with bell-shaped profile **54** and surrounding it an optically inactive cone joining the upper surface, with straight-line profile **55** and pegs **56** going into circuit board **57**. Illuminator **50** further comprises LED package **58** with emissive chip **58C** immersed in transparent hemispheric dome **58D**. The term 'toroid' distinguishes from the conventional term 'torus', which solely covers the case of zero tilt angle. The highly oblique lighting setup of refrigerator-cabinet shelf-fronts

quite understandably involves tilting the torus so that the lensing effect of the elliptical arc points toward the center of the shelf.

Arc **52** of FIG. **5** extends to tilt angle τ , which in this case is 17° , its importance being that it is the tilt angle of major axis **52A** of elliptical arc **52**. Its minor axis **52B** lines up with the radius at the edge of arc **51**, ensuring profile-alignment with equal surface tangency. There are three free parameters which define a particular outer surface of illumination lens **51**, as intended for different shelf distances. The first is the radius of arc **52**, which controls the amount of de-magnification by the central portion of illumination lens **51**. The second is the tilt angle τ , which defines the orientation of elliptical arc **53**, namely towards the shelf center of FIG. **1**. The third free parameter of the upper surface is the ratio of major axis **53A** to minor axis **53B**, in this case 1.3:1, defining the above-discussed source magnification. Ray-fan **59** comprises central rays (i.e., originating from the center of chip **58C**) at 2° intervals of off-axis angle. The central ten rays designated by dotted arc **59C** illustrate the diverging character of the center of lens **51**, which provide the central demagnification required for uniform illumination. The remaining rays are all sent at steep angles to the horizontal, providing the lateral source magnification of FIG. **4**.

The central cavity surrounding LED **58** has bell-shaped profile **54** defined by the standard aspheric formula for a parabola (i.e., conic constant of -1) with vertex at z_v , vertex radius of curvature r_c , 4^{th} -order coefficient d and 6^{th} -order coefficient e :

$$z(x)=z_v+x/r_c+dx^4+ex^6$$

In order for profile **54** to arc downward rather than upward, the radius of curvature r_c is negative. The aspheric coefficients provide an upward curl at the bottom of the bell, to help with cutting off the illumination pattern. The particular preferred embodiment of FIG. **5**, with a cavity entrance-diameter set at 6.45 mm, is defined by:

$$z_v=6 \text{ mm } r_c=-1.69 \text{ mm } d=-0.05215 \text{ } e=0.003034$$

This profile only needs minor modification to be suitable for preferred embodiments illuminating other shelf distances.

FIG. **6A** through **6F** shows illumination lens **60** and LED chip **61**. In FIG. **6A**, rays **62** come from points on the shelf at the indicated x coordinates of 0, 2", and 4" laterally from the lens. Each bundle is just wide enough that its rays end at the edges of chip **61**, which is the definition of a source image. Each bundle is narrower than chip **61** would appear by itself, in accordance with the previously discussed demagnification. The central portion of lens **60** that is traversed by rays **62** can be seen to be a concave, diverging lens, as previously mentioned.

FIG. **6B** shows ray bundle **63** proceeding from the distance x_m to the maximum of the illumination pattern in FIG. **2**. It is twice the width of those in FIG. **6A**.

FIG. **6C** shows ray bundle **64** proceeding from the distance x_m to the middle of the shelf, as shown in FIG. **2**.

FIG. **6D** shows ray bundle **65** proceeding from beyond mid shelf, at 18".

FIG. **6E** shows ray bundle **66** proceeding from beyond mid shelf, at 20", nearly filling the lens. This is the maximum source magnification this sized lens can handle.

FIG. **6F** shows ray bundle **67** proceeding from the edge of the illumination pattern, at $x_E=22"$. Note that these rays miss chip **61**, indicating that there will be no light falling there, which is required by the pattern cutoff.

The progression of FIG. **6A** through **6F** is the basis for the numerical generation of the upper and lower surface profiles

of the lens, starting at the center and working outwards, as will be disclosed below. The results of this method can sometimes be closely approximated by the geometry of FIG. **5**.

The illumination lens of FIG. **5** has elliptical and aspheric-parabolic surfaces with shapes that are exactly replicable by anyone skilled in the art. In the illumination pattern of FIG. **2**, the central depression to $\frac{3}{4}$ the maximum value was empirically found to work with the lens array of FIG. **1**, with six lenses on each side. This lens is the first commercially available design enabling only six LEDs to be used, rather than the dozen or more of the prior art. More recently, however, even higher-power LEDs have become available that only require two per door, as FIG. **7** illustrates.

FIG. **7** shows rectangular outline **70** representing a typical refrigerator door that is 30" wide and 60" high, with other doors, not shown, to either side. Dashed rectangles **71** denote the mullions behind which the shelf lighting is mounted, typically at 3-6" from the front of the illuminated shelves. This is much closer than the distance to the shelf center, denoted by centerline **72**. There are four illuminators (two on either side), denoted by small circles **73**. Each illuminator produces an illuminated circle with its peak on a ring denoted by solid circles **74** and its edge on dotted circles **75**. Here the circles **74** have radius of about a fifth of the shelf width, or a third the way to centerline **72**. The circles **75**, where illumination has fallen to zero, are sized to reach nearly all the way across the shelf. As in FIG. **1**, each pattern has the value $\frac{1}{2}$ at centerline **72**, so two lenses add to unity. Also, at shelf center-point **76** the four patterns overlap, so at this distance each pattern must have the value $\frac{1}{4}$, and thus add to unity. This same configuration is applicable for LCD backlights comprising square-arrayed LEDs, merely on a smaller scale. This arrangement of precisely configured illumination lenses is capable of generating uniformity satisfactory for LCD backlights.

The LEDs used in the arrangement of FIG. **7** must be three times as powerful as those used for FIG. **1**. This greater flux has unwanted consequences of triply enhanced scattered light, strengthened even more by the greater size of the lenses used for FIG. **7** versus the smaller ones which would suffice for FIG. **1**. The illumination pattern of FIG. **2** has a central dip in order to compensate for the close spacing of the lenses. When scattering is significant, however, the scattered light can be strong enough to provide all the illumination near the lens. The upshot is that the illumination pattern shown in FIG. **2** would have nearly zero intensity on-axis. The resultant lens has a previously unseen feature: either or both surfaces have a central cusp that leaves no direct light on the axis, resulting in a dark center for the pattern, in order to compensate for the scattered light.

FIG. **8** is a cross-section of illuminator **80**, comprising circularly symmetric illumination lens **81**, bounded by an upper surface comprising a central cusp **82** formed by a surrounding toroid with tailored arc **83** as its profile. Lens **81** is also bounded by a lower surface comprising a central cavity with tailored profile **84** preferably peaking at its tip, and surrounding it an optically inactive cone joining the upper surface, with straight-line profile **85** and pegs **86** going into circuit board **87**. Illuminator **80** further comprises centrally located LED package **88** with emissive chip **88C** immersed in transparent hemispheric dome **88D**.

The optically active profiles **83** and **84** of FIG. **8** are said to be tailored due to the specific numerical method of generating it from an illumination pattern analogous to that of FIG. **2**, but with little or no on-axis output. The reason for this is, as aforementioned, to compensate for real-world scattering from the lens. The profiles **83** and **84** only control light propa-

gating directly from chip **83C**, through dome **83D**, and thence refracted to a final direction that ensures attainment of the required illumination pattern. This direct pattern will be added to the scattering pattern of indirect light, which thus needs to be determined first.

FIG. **9** shows illumination lens **91**, identical to lens **81** of FIG. **8**, with other items thereof omitted for clarity. From LED chip **98C** issues ray bundle **92**, comprising a left ray (dash-dot line), a central ray (solid line), and a right ray (dashed line), issuing respectively from the left edge, center, and right edge of LED chip **98C**. Anywhere within lens **91**, these rays define the apparent size of chip **98C** and thus how much light is passing through a particular point. Any light scattered from such a point will be a fixed fraction of that propagating light. The closer to the LED the more light is present at any point, and the greater the amount scattered. This scattering gives the lens its own glow, separate from the brightness of the LED itself when directly viewed.

Strictly speaking, scattering does not take place at a point but within a small test volume, shown as infinitesimal cube **93** in FIG. **9**, magnified for clarity. It is oriented along the propagation direction of ray bundle **92**. It has cross-section **93A** of area dA and propagation length dl , such that its volume is simply $dV=dl dA$. Within cube **93** can be seen the left, central, and right rays of bundle **92**, now switched sides. The right and left rays define solid angle Ω , indicating the apparent angular size of LED chip **98C** as seen from cube **93** within lens **91**. The greater this solid angle the more light will be going through cube **93**. LED chip **98C** has luminance L , specified in millions of candela per square meter. This is reduced when ray bundle **92** goes into lens **91**, due to less-than-unity transmittance τ caused by Fresnel reflections. Going into cube **93** the ray bundle **92** has intensity I given simply by $I=\tau L dA$. The total flux F passing through cube **93** is then given, simply again, by $F=I\Omega$.

Volume scattering removes a fixed fraction of this intensity I per unit length of

$$I(l)=I(0)e^{-\kappa l}$$

propagation, similar to absorption. Both are described by Beer's law:

Here $I(0)$ is the original intensity and $I(l)$ is what remains after propagation by a distance l , while scattering coefficient κ has the dimension of inverse length. It can easily be determined by measuring the loss in chip luminance as seen through the lens along the path l of FIG. **9**.

Returning to cube **93v** of FIG. **9**, the ingoing intensity I is reduced by the small amount $dI=e^{-\kappa dl}$. This results in a flux decrement $dF=dI\Omega$ that is subtracted from F . Then the emission per unit volume is dF/dV . Integrating this over the entire optically active volume of lens **91** gives the total scattered light. FIG. **9** further shows observer **94** gazing along line of sight **95**, along which direct rays **97** give rise to scattering points **96**, summing into a lens glow that acts as a secondary light source surrounding the LED.

These scattering phenomena are usually looked upon as disadvantageously parasitic, acting only to detract from optical performance. There is a new aspect to this, however, where some volume scattering would be beneficial. It arises in the subtle failings of current high-brightness LEDs, namely that of not delivering the same color in all directions. More specifically, many commercially available LEDs with multi-hundreds of lumens output, look much yellower when seen laterally than face-on. This is because of the longer path through the phosphor taken by light from the blue chip.

Thick phosphors have uniform whiteness, or color temperature, in all directions, but they reduce luminance due to

the white light being emitted from a much bigger area than that of the blue chip. Conformal coatings, however, are thin precisely in order to avoid enlarging the emitter, but they will therefore scatter light much less than a thick phosphor and therefore do much less color mixing. As a result, lateral light is much yellower (2000 degrees color temp) and the face-on light much bluer (7000 degrees) than the mean of all directions. As a result of this unfortunate side-effect of higher lumen output, the lenses disclosed herein will exhibit distinct yellowing of the lateral illumination, and a distinct bluing of the vertical illumination.

The remedy for this inherent color defect is to use a small quantity of blue dye in the lens material. Since the yellow light goes through the thickest part of the lens, the dye will automatically have its strongest action precisely for the yellowest of the LEDs rays, those with larger slant angles. The dye embedded in the injection-molding material should have an absorption spectrum that only absorbs wavelengths longer than about 500 nm, the typical spectral crossover between the blue LED and the yellow phosphor. The exact concentration will be inversely proportional to lens size as well as to the absorption strength of the specific dye utilized.

A further form of scattering arises from Fresnel reflections, aforementioned as reducing the luminance of rays as they are being refracted. FIG. **9** further shows first Fresnel-reflected ray **92F1** coming off the inside surface of lens **91**, then proceeding into the lens to be doubly reflected out of the lens onto circuit board **92F1**. This ray has strength of $(1-\tau)$ relative to the original ray **92**. Of similar strength is the other Fresnel-reflected ray, **92F2**, which proceeds from the outer surface to the bottom of the lens. These two rays are illustrative of the general problem of stray light going where it isn't wanted. Unlike the volume scattering at points **96**, these Fresnel-reflected rays can travel afar to produce displeasing artifacts. It has been well-known for many decades of optical engineering that the easiest way to deal with this is to institute surface scattering of these stray rays. Since the flat conical bottom surface **91C** of lens **91** intercepts most of these stray Fresnel reflections, the tried-and-true traditional solution is simply to roughen the corresponding mold surface so that the Fresnel light is dissipated to become part of the above-described volume scattering. At the termination of ray **92F2** can be seen the scattered rays, some of which illuminate the top of board **99**, which of course would also be scattering, such as by white paint.

FIG. **10** shows graph **100** with abscissa **101** denoting distance in millimeters from the center of the lens of FIG. **9** and ordinate **102** denoting illuminance relative to the pattern maximum (in order to generalize to any illumination level). Dashed curve **103** is the ideal illumination pattern desired for the configuration of FIG. **7**, given an inter-lens spacing of 125 mm and a target distance of 23 mm. These dimensions represent a backlight application, where the LEDs are arrayed within a white-painted box, and the target is a diffuser screen, with a liquid-crystal display (LCD) just above it. Increased LED luminosity mandates fewer LEDs, to save on cost, while aesthetics push for a thinner backlight. These two factors comprise a design-pressure towards very short-throw lighting.

The 'conical pattern' of curve **103** and its converse (not shown) from an illuminator at 125 mm, will add to unity, which assures uniform illumination. Dash-dot curve **104** depicts the combined parasitic illuminance on that target plane caused by the above-discussed volume and surface scattering from a lens at $x=0$. This curve is basically the cosine⁴ of the off-axis angle to a point on the target. Solid curve **105** is the normalized difference between the other two

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curves, representing the pattern that when scaled will add to curve **104** to get a total illuminance following curve **103**. In this case the scattered light of curve **104** is strong enough to deliver 100% of the required illuminance just above the lens. In such a case the central cusp **82** of FIG. **8** will ensure that the central illuminance is zero when only counting direct light that is delivered through the lens.

The illumination pattern represented by curve **105** of FIG. **10** can be used to numerically generate the inner and outer profiles of the lens **81** of FIG. **8**, utilizing rays from the right and left edges of the source. Dotted curve **106** of FIG. **10** graphs the relative size of the source image height (as shown in FIG. **6A-F**) required by the illuminance pattern of curve **105**. This height function is directly used to generate the lens profiles.

FIG. **11** shows LED **110** and illumination lens **111**, of 20 mm diameter, sending right ray **112** and left ray **113** to point **114**, which has coordinate x on planar target **115**, located 23 mm above LED **110**. Right ray **112** hits point **114** at slant angle γ , and left ray **113** at slant angle $\gamma + \Delta\gamma$. In the two-dimensional analysis of FIG. **11**, the illuminance $I(x)$ at point x is proportional to the difference between the sines of the left and right rays' slant angles:

$$I(x) \propto \sin(\gamma + \Delta\gamma) - \sin(\gamma)$$

This angular requirement can be met by the proper height H of the source image, namely the perpendicular spacing between right ray **112** and left ray **113**, at the lens exit of **112**. Curve **106** of FIG. **10** is a plot of this height H , relative to its maximum value. From this geometric requirement the lens profiles can be directly generated by an iterative procedure that adds new surface to the previously generated surface.

FIG. **12** shows incomplete illumination lens **121**, positioned over LED **120**. It is incomplete in that it represents a typical iteration-stage of generating the entire lens of FIG. **11**. The portion of Lens **111** of FIG. **11** that is shown as a slightly thickened curve terminates at its intersection, shown as point **124**, with right ray **122**. In FIG. **12**, a new left ray **123** is launched that is barely to the right of left ray **113** of FIG. **11**. After going through terminal point **126** and then through previously generated upper surface **121**, it will intercept the target (not shown) at a new point $x + dx$, just to the right of point x of FIG. **10**. This point will have an already calculated source-height requirement such as curve **105** of FIG. **10**, fulfilled by launching a new right ray **122** from $x + dx$. Ray **122** will intercept the lens surface at point **125**, upon new surface that has been extended from point **124**. The new surface has a slope determined by the necessity to deflect ray **122** towards point **126** on the interior surface of lens **121**. The location of this point **127** is determined by right ray **122-S** coming from the right edge of LED chip **120C**. The off-axis angle of this ray is determined by the usual requirement that the angular-cumulative intensity of right ray **122S** equal the spatially cumulative illumination at point $x + dx$, which is known from the desired illumination pattern, such as that shown by curve **105** of FIG. **10**. The slope of this new interior surface, from point **126** to new point **127**, is determined by the necessity of refracting ray **122S** so it joins ray **122** to produce the proper source-image height for the illumination of the target at point $x + dx$. In this fashion, the generation of lens **121** will be continued until all rays from chip **120C** are sent to their proper target coordinates, and its full shape is completed.

The profile-generation method just described is two-dimensional and thus does not account for skew rays (i.e., out-of-plane rays), which in the case of a relatively large source can give rise to noticeable secondary errors in the output pattern, due to lateral variations in the size of the

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source image. This effect necessitates a fully three-dimensional source-image analysis for generating the lens shape, as shown in FIG. **13**.

The lens-generation method of FIG. **12** traces left ray **123** through the previously generated inner and outer surfaces to a target point with lateral coordinate $x + dx$. The pertinent variable is the height H of the source image. In three dimensions, however, rays must be traced from the entire periphery of the LED's emission window out to the target point, where they limit the image of the source as seen through the lens from that point. An illumination lens acts to alter the sources' apparent size from what it would be by itself. The size of the source image is what determines how much illumination the lens will produce at any target point.

FIG. **13A** is a schematic view from above of circular illumination lens **130**, with dotted lines showing is incomplete, its design iteration having only extended so far to boundary **131**. Circular source **132** is shown at the center of lens **130**, and oval **133** represents the source image it projects to target point $x + dx$ (not shown). This source image is established by reverse ray tracing from the target point back through the lens to the periphery of the source. The source image is the oval outline **133** on the upper surface where these rays intercept it. Thus the already completed part of the lens will partially illuminate the target point, and a small element of new surface must be synthesized for full illumination.

FIG. **13B** is a close-up view showing source ellipse **133** and boundary **131**, also showing curve **134**, representing a small element of new surface that will be added in order to complete source image **133** and achieve the desired illumination level at target point $x + dx$. Of course, when new upper surface is added there will have to be a corresponding element of new lower surface added as well. Just as enough there must be enough new upper surface to finish the source image, so too must there be enough lower surface to provide the source image to the upper surface. Since both the extent and slope of this new lower surface must be determined as free variables, the design method must be able to calculate both unknowns, but in general the slopes of the new elements of upper and lower surface will be smooth continuations of the previous curvatures of the surfaces.

Traditionally, non-imaging optics deals only with rays from the edge of the source, but the illumination lenses disclosed herein go beyond this when assessing the source image at each target point. The incomplete source image of FIG. **13B** gives rise to a less-than-required illuminance at the target point of interest, at lateral coordinate x . In order to calculate this illuminance, however, rays must be reverse traced back to the entire source, not just its periphery. This is especially true when the source has variations in luminance and chrominance. Then the flux from each small elemental area dA of FIG. **13B** is separately calculated and integrated over the source image as seen through already completed surface. The deficit from the required illuminance will then be made up by the new surface **134** of FIG. **13B**. Its size is such that the additional source image area will just finish the deficit. When the illumination pattern only changes gradually, as with the linear ramps just discussed, the deficit is always modest because the previously generated surface has done a good job of getting close to the required illuminance. The new surface will not have to scrunch the new source image, due to a tiny deficit, nor expand it wildly for a large deficit, because the target pattern is 'tame' enough to prevent this.

This design method can be called 'photometric non-imaging optics', because of its utilization of photometric flux accounting in conjunction with reverse ray tracing to augment the edge-ray theorem of traditional non-imaging optics.

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The iterative process that numerically calculates the shape of a particular illumination lens can begin, alternatively, at either the center or the periphery. If the lens diameter is constrained, the initial conditions would be the positions of the outer edges of the top and bottom surfaces, which then totally determines the lens shape, in particular its central thickness. If this thickness goes below a minimum value then the initial starting points must be altered. While this is conceptually feasible, in practical terms it leaves the problem underdetermined, whereas the reverse ray tracing of FIG. 13A utilizes the previously generated surface via reverse ray tracing. Thus it is easier to begin the design iteration at the center of the lens using some minimum thickness criterion, e.g., 0.75 mm. The height of the lens center above the source would be the primary parameter in determining the overall size of the lens. The other prime factor is how the central part of the lens is configured as a negative lens, that is, whether concave-plano, concave-concave, or plano-concave. Also, a concave surface can either be smooth or have the cusp-type center as shown in FIG. 8, in the case of strong parasitic losses.

FIG. 14 shows concave-plano lens-center 140, to be used as a seed-nucleus for generating an entire illumination lens. Its diameter is determined by the width of ray fan 141, which propagates leftward from the target center (not shown) $x=0$ at a distance of 23 mm above (to the right of) LED chip 142, the size of which has been exaggerated for clarity. Ray fan 141 has the width necessary to achieve the desired illumination level at the center of the target, and in short-throw lighting this is less than what the LED would do by itself. This means the central part of the illumination lens must demagnify the source, which is why the lens-center is diverging, with negative focal length. In fact, the very function of lens-center 140 is to provide the proper size of source image (of which ray fan 141 is a cross-section) for the target center, $x=0$.

FIG. 14 also shows expanding ray fan 143, originating at the left edge of chip 142. This will mark the upper edge of a source image as seen from the x -positions at which these left rays intercept the target plane (not shown, but to the right). These rays exemplify how edge rays are sent through previously established surfaces.

FIG. 15 shows concave-concave lens-center 150, central ray-fan 151, chip 152, and left-ray fan 153. The lens surfaces have about half the curvature of the concave surface of FIG. 14.

FIG. 16 shows plano-concave lens-center 160, central ray-fan 161, chip 162, and left-ray fan 163.

In the progression from FIGS. 14 to 16, the lowest left ray (the one with an arrow) lies at a shrinking slant angle γ r , indicating different illumination behavior and setting a different course towards the final design. All three configurations produce the same illuminance at target $x=0$, that is to say the same size source image, as shown by the ray fans 141, 151, & 161 being of identical size as they arrive at each lens, which is equivalent to saying they produce the same target illuminance at center $x=0$.

In FIG. 16, a plano-concave lens-center 160 is defined by a ray-fan 161, which proceeds from the center of the short-throw target (not shown). The width of the ray-fan 161 as it enters the lens-center 160 is determined by an illumination prescription, which may be less than what LED chip 162 would illuminate by itself. Hence ray-fan 161 is seen to be expanded by the lens-center 160 so it diverges to cover LED chip 162.

The diverging ray-fan 163 proceeds from the left edge of the LED chip 162, and goes through the lens-center 160 and then out to the short-throw target. These rays from the ray-fan

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163 help guide the synthesis of the rest of the illumination lens, as shown in FIGS. 19 and 20 below.

In FIG. 16, ray-fan 163 diverges more for this type of lens-center than those of the FIGS. 14 and 15, as shown by smaller slant angle γ r . Thus the plano-concave type of illumination lens shown in FIGS. 19 and 20 may be appropriate for illumination prescriptions involving arrays of lenses rather than just one acting alone, for instance.

FIG. 17 shows illumination lens 170, numerically generated from a concave-plano center-lens, as in FIG. 14. Planar source 171 is the light source from which it was designed.

FIG. 18 shows illumination lens 180, numerically generated from a concave-concave center-lens, as in FIG. 15. Planar source 181 is the light source from which it was designed.

FIG. 19 shows illumination lens 190, numerically generated from a plano-concave center-lens, as in FIG. 16. Planar source 191 is the light source from which it was designed.

FIG. 20 shows a short throw (i.e. illumination) lens 200, and its related geometry. FIG. 20 is similar to the illumination lens of FIG. 19, but is shown in more detail. The short throw lens 200 and the geometry thereof may be further explained as described above with respect to FIGS. 3-5. The short throw lens 200 may be useful as a BLU lens for LCD displays.

A planar light source 201, which may be a light emitting diode, is arranged under the short throw lens 200. The short throw lens 200 has inside and outside surfaces 204 and 206, respectively, while 208 is the bottom surface of the short throw lens 200 connecting the inside and outside surfaces 204 and 206. A central cavity 202 of the short throw lens 200 comprises a flat top 203 and lateral flanks 204. The flat top 203 is substantially parallel to an extension direction of the light source 201. The outer surface of the short throw lens 200 comprises a central indentation 205 and surrounding toroid 206. The central indentation 205 acts as a diverging lens, as described above with respect to FIG. 16.

Tilted axis 207 aims in the direction of maximum source-image size, towards the place on a short-throw target (e.g., an external surface to be illuminated, not shown), where the short throw lens 200 has the greatest magnification. The exact location of this place depends upon the particular illumination prescription, but it is generally about halfway to the edge of the target. Thus, the tilted axis 207 is a line drawn from the center point 0 of the light source 201 through the point of maximum magnification on the second surface 206, which points to the midpoint of the short-throw target.

The bottom surface 208 of the short throw lens 200 is inclined by an angle τ_1 , which represents the angle that that the bottom surface 208 makes with respect to the horizontal optical axis X . Angle τ_1 creates the tilted axis 207. The short throw lens 200 may be ellipsoidal, and the bottom surface 208 may be inclined by the angle τ_1 along the major and minor axes of the ellipsoid.

R_1 is the radius vector representative of a light ray leaving the light source 201 and incident on the flat top 203 or the inside surface 204 of the central cavity 202 of at point P1. Angle α_1 represents the angle that that R_1 makes with respect to the vertical optical axis Z . It should be noted that the central cavity 202 is axisymmetrical with respect to optical axis Z . Radius vector R_1 has monotonically increasing values, represented by ΔR_1 in the range of 0 to 25° as α_1 increases incrementally ($\Delta\alpha_1$). This is due to the fact that the flat top 203 of the central cavity 202 is substantially flat in the angular range of 0 to 25°. FIG. 21 shows a detailed view of these features.

R_2 is the radius vector to the light ray exiting the outside surface 206 of the short throw lens 200 at point P2, while α_2 represents the angle that that ray makes with respect to the

vertical optical axis Z, and ΔR_2 is the incremental change in R_2 as α_2 increases incrementally ($\Delta\alpha_2$). R_2 monotonically increases as α_2 increases in a range of 0 to 60°. The toroidal portion of the outside surface **206** is axisymmetrical with respect to the vertical optical axis Z. In order to avoid total internal reflection of light emitted from the light source **201** into the short throw lens **200**, $A_2 < 1/\sqrt{(n^2-1)}$, where n is a refraction index of a material forming the illumination lens, and $A_2\Delta R_2/(R_2\Delta\alpha_2)$. Here, the short throw lens **200** is made of a transparent material having an index of refraction between 1.41 and 1.65.

The painted bottom surface **208** of the short throw lens **200** may be configured to absorb internal (e.g., Fresnel) reflection of light emitted by the light source **201**. The bottom surface **208** may be painted such that the absorptivity of the bottom surface **208** is at least 80%, and preferably 90%. The bottom surface **208** of the short throw lens **200** may be painted with a black paint in order to increase absorptivity. Alternatively, the top of a printed circuit board on which the short throw lens **200** may be mounted (not shown), can be similarly painted in order to increase absorptivity.

These three lenses were designed utilizing rays from the periphery of the light source, in this case circular. The size of the lens is a free parameter, but etendue considerations dictate that a price be paid for a lens that is too small. In the case of a collimator, the output beam will be inescapably wider than the goal if the lens is too small. In the case of the short-throw illumination lenses disclosed herein, the result will be an inability to maintain an output illumination pattern that is the ideal linear ramp of curve **103** of FIG. **10**, because it requires the source image of curve **106**. If the lens is smaller than the required source image size, then it cannot supply the required illumination. Thus the lens size will be a parameter fixed by the goal of a linear ramp. Lenses that are too small will have some rays trapped by total internal reflection instead of going to the edge of the pattern. If this is encountered in the design process then the iteration will have to re-start with a greater height of the lens-center above the LED.

LED illumination lenses may need a circularly symmetric illumination lens that can be used in small numbers (such as five or six per mullion) and still attain uniformity, because the individual patterns are such that those few may add up to caustic-free uniformity. Thus, an objective of the present invention is to provide a lens with a circular illumination pattern that multiples of which will add up to uniformity across a rectangle. It is a further objective of the present invention to attain a smaller lens size than the above mentioned approaches, leading to device compactness that results in lower manufacturing cost.

The smaller lens size can be achieved by a specific tailoring of its individual illumination pattern. This pattern is an optimal annulus with a specific fall-off that enables the twelve patterns to add up to uniformity between the two illuminating mullions upon which each row of six illuminators are mounted. This fall-off at the most oblique directions is important, because this is what determines overall lens size. The alternative approaches are: (1) Each mullion illuminates 100% to mid-shelf and zero beyond, which leads to the aforementioned caustics; (2) Each mullion contributes 50% at the mid-point, falling off beyond it. The latter is the approach of the present invention.

In conclusion, the exemplary embodiments disclosed herein fulfill a most challenging illumination task, the uniform illumination of close planar targets by widely spaced lenses. Deviations from this lens shape that are not visible to casual inspection may nevertheless suffice to produce detractive visual artifacts in the output pattern. Experienced mold-

ers know that sometimes it is necessary to measure the shape of the lenses to a nearly microscopic degree, so as to adjust the mold-parameters until the proper shape is achieved.

Experienced manufacturers also know that LED placement is critical to illumination success, with small tolerance for positional error. Thus a complete specification of a lens shape necessarily requires a high-resolution numerical listing of points mathematically generated by a fully disclosed algorithm. Qualitative shape descriptors mean nothing to computer-machined injection molds, nor to the light passing through the lens. Unlike the era of manual grinding of lenses, the exactitude of LED illumination lenses means that without a numerical method of producing those lens-profile coordinates, there will be no lens.

It will be apparent to those skilled in the art that various modifications and variation can be made in the present invention without departing from the spirit or scope of the invention. Thus, it is intended that the present invention cover the modifications and variations of this invention provided they come within the scope of the appended claims and their equivalents.

What is claimed is:

1. A light-emitting device, comprising:

a light-emitting element; and

an illumination lens configured to control light emitted from the light-emitting element, the illumination lens comprising a first surface from which light emitted from the light-emitting element enters the illumination lens, wherein the first surface comprises a central cavity that is axisymmetrical with respect to a reference axis of the light-emitting device,

wherein R_1 monotonically increases as α_1 increases in a range of $\alpha_1 < 25^\circ$, where an intersection of the reference axis and an emission surface of the light-emitting element is a reference point, R_1 is a distance between the reference point and an arbitrary point on the first surface, and α_1 is an angle between the reference axis and a straight line that passes through the reference point and the arbitrary point on the first surface.

2. The light-emitting device of claim 1, wherein the illumination lens further comprises a second surface from which the light being incident on the first surface is emitted from the illumination lens,

wherein the second surface comprises a toroidal portion that is axisymmetrical with respect to the reference axis and a central indentation comprising a portion of the second surface surrounded by the toroidal portion, the central indentation encompassing an intersection of the reference axis and the second surface, the central indentation acting as a diverging lens.

3. The light-emitting device of claim 2, wherein R_2 monotonically increases as α_2 increases in a range of $\alpha_2 < 60^\circ$, where R_2 is a distance between the reference point and an arbitrary point on the second surface, and α_2 is an angle between the reference axis and a straight line that passes through the reference point and the arbitrary point on the second surface.

4. The light-emitting device of claim 3, wherein $A_2 < 1/\sqrt{(n^2-1)}$, where n is a refraction index of a material forming the illumination lens, and $A_2 = \Delta R_2/(R_2\Delta\alpha_2)$, where ΔR_2 is an increment of R_2 and $\Delta\alpha_2$ is an increment of α_2 .

5. The light-emitting device of claim 1, wherein the first surface central cavity comprises a flat top, the flat top being substantially parallel to an extension direction of the light-emitting element.

6. The light emitting device of claim 1, wherein the illumination lens comprises a transparent material having an index of refraction between 1.41 and 1.65.

7. The light emitting device of claim 1, wherein the light emitting element comprises a light emitting diode (LED). 5

8. The light emitting device of claim 2, wherein the bottom surface of a conical section of the second surface toroidal portion is inclined by an angle τ .

9. The light emitting device of claim 8, wherein a line drawn from the emission surface of the light emitting element through the point of maximum magnification on the second surface toroidal portion points to the midpoint of an external surface to be illuminated. 10

10. The light emitting device of claim 8, wherein the illumination lens is ellipsoidal, and the bottom surface of the conical section of the second surface toroidal portion is inclined by the angle τ along the major and minor axes thereof. 15

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